

Critical RSOS and Minimal Models II: Building Representations of the Virasoro Algebra and Fields

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Abstract

We consider $sl(2)$ minimal conformal field theories and the dual parafermion models. Guided by results for the critical A_L Restricted Solid-on-Solid (RSOS) models and its Virasoro modules expressed in terms of paths, we propose a general level-by-level algorithm to build matrix representations of the Virasoro generators and chiral vertex operators (CVOs). We implement our scheme for the critical Ising, tricritical Ising, 3-state Potts and Yang-Lee theories on a cylinder and confirm that it is consistent with the known two-point functions for the CVOs and energy-momentum tensor. Our algorithm employs a distinguished basis which we call the L_1 -basis. We relate the states of this canonical basis level-by-level to orthonormalized Virasoro states.

1 Introduction

In this series of papers we study the relation between the critical Restricted Solid-on-Solid (RSOS) lattice models [1] and the minimal conformal field theories (CFT) [2, 3]. In Part I of this series (FPI) [4] we argued that, for these theories on a cylinder, the Hilbert space of physical states consists of fermionic states built from one-dimensional configurational RSOS paths on the lattice. As this series of papers will show, this allows us to translate much of the insight gained from studying finite-size critical RSOS lattice models by Yang-Baxter techniques [5] across to the study of the associated CFTs.

For some time now, workers in the field have been trying to build matrix representations of the Virasoro algebra based on paths [6, 7] and to make sense of a finitized Virasoro algebra [8, 9]. In this paper, we propose a general algorithm to build matrix representations of the Virasoro generators and chiral vertex operators (CVOs) and implement it for the critical Ising, tricritical Ising, 3-state Potts and Yang-Lee theories on a cylinder. We emphasize that our algorithm constructs these matrices level-by-level and, while our algorithm appears to be consistent level-by-level, we do not have expressions for these matrices valid at arbitrary level.

The layout of the paper is as follows. In Section 2, we recall some relevant facts about CFT, the Virasoro algebra, Virasoro states, primary fields and their two-point functions. In Section 3, we present our general algorithm to build matrix representations of the Virasoro generators and primary fields. We implement this algorithm for the critical Ising, tricritical Ising, 3-state Potts and Yang-Lee theories on a cylinder in Sections 4–7. We conclude in Section 8 with a discussion of open questions for further research.

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2 Conformal Field Theory

2.1 Minimal models and \mathbb{Z}_k parafermions

The $sl(2)$ minimal models [2] $\mathcal{M}(p', p)$ with p, p' coprime have central charges

$$c = 1 - \frac{6(p-p')^2}{pp'} \quad (2.1)$$

The conformal weights are

$$h = h_{r,s} = \frac{(rp - sp')^2 - (p - p')^2}{4pp'}, \quad r = 1, 2, \dots, p' - 1; \quad s = 1, 2, \dots, p - 1 \quad (2.2)$$

and the Virasoro characters are

$$\chi_h(q) = \frac{q^{-c/24+h}}{(q)_\infty} \sum_{k=-\infty}^{\infty} (q^{k(kpp'+rp-sp')} - q^{(kp+s)(kp'+r)}) \quad (2.3)$$

where

$$(q)_n = \prod_{k=1}^n (1 - q^k) \quad (2.4)$$

The minimal models are unitary if $p - p' = \pm 1$ and non-unitary otherwise. We consider only the diagonal A -type series with $p' < p$ and use the critical Ising $\mathcal{M}(3, 4)$, tricritical Ising $\mathcal{M}(4, 5)$ and Yang-Lee theories $\mathcal{M}(2, 5)$ as prototypical examples.

The $sl(2)$ \mathbb{Z}_k parafermion models [10] are dual to the minimal models and have central charges

$$c = \frac{2(k-1)}{k+2}, \quad k = 2, 3, \dots \quad (2.5)$$

We consider only the diagonal A -type series and we use the \mathbb{Z}_3 or hard hexagon model [11, 5] as the prototypical example. The hard hexagon model is in the universality class of the 3-state Potts model so we refer to this as the 3-state Potts CFT. Generally, the characters of the \mathbb{Z}_k models are string functions but, for the \mathbb{Z}_3 model, these are easily related to the Virasoro characters of the $\mathcal{M}(5, 6)$ model.

The minimal and \mathbb{Z}_k parafermion models are rational and admit a finite number of primary fields $\phi(z) = \phi^{(h)}(z)$. We view these theories as arising from the continuum scaling limit of the A_L RSOS models [1] with $L = p - 1$ and $L = k + 1$ respectively.

2.2 Virasoro algebra and Virasoro states

The Virasoro algebra

$$\text{Vir} = \langle L_n, n \in \mathbb{Z} \rangle \quad (2.6)$$

is an infinite dimensional complex Lie algebra associated with conformal symmetry. The generators L_n satisfy the commutation relations

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12} n(n^2 - 1)\delta_{n,-m} \quad (2.7)$$

where the central element c is the central charge. These relations remain unchanged under an orthogonal change of basis $L_n \mapsto L'_n = U^T L_n U$. The Virasoro generators are the modes of the energy-momentum tensor

$$T(z) = \sum_{n \in \mathbb{Z}} L_n z^{-n-2} \quad (2.8)$$

On a cylinder with prescribed boundary conditions, which is the case of primary concern here, there is just one copy of the Virasoro algebra. For bulk theories on the torus, however, there is a second copy $\overline{\text{Vir}}$ of Virasoro which is the antiholomorphic counterpart.

For rational CFTs, the Hilbert space \mathcal{H} of states on which Vir acts is naturally decomposed into a finite direct sum of irreducible highest weight representations (Virasoro modules)

$$\mathcal{H} = \bigoplus_h \mathcal{V}_h \quad (2.9)$$

where the sum is over the conformal weights h of the primary fields $\phi(z) = \phi^{(h)}(z)$. The vacuum $|0\rangle$ and primary (highest weight) states $|h\rangle$ are characterised by

$$L_0|h\rangle = h|h\rangle; \quad L_n|0\rangle = 0, \quad n \geq -1; \quad L_n|h\rangle = 0, \quad n > 0 \quad (2.10)$$

Moreover, there is a one-to-one correspondence between primary fields and primary states induced by

$$\lim_{z \rightarrow 0} \phi^{(h)}(z)|0\rangle = |h\rangle \quad (2.11)$$

The vacuum state $|0\rangle$ with $h = 0$ corresponds to the identity operator.

The generically reducible highest weight representation of Vir (Verma module) is the linear span of Virasoro states in the canonical form

$$L_{-n_j} L_{-n_{j-1}} \dots L_{-n_1} |h\rangle, \quad n_j \geq n_{j-1} \geq \dots \geq n_1 \geq 1 \quad (2.12)$$

If its maximal proper submodule is quotiented out, we are led to the irreducible Virasoro module $\mathcal{V}_h = \mathcal{V}_{c,h}$ and the states (2.12) are no longer linearly independent due to the existence of null vectors. The generic Virasoro module in the $h = 0$ vacuum sector is shown in Figure 1. Typically, for given c and h , some states at a given level enter in a vanishing non-trivial linear combination that is the null vector. Surprisingly, it seems that a complete set of linearly independent Virasoro states is not known even for the Ising model, although it is known [6] for the Yang-Lee theory $\mathcal{M}(2, 5)$ and the whole family $\mathcal{M}(2, 2n + 3)$, $n \geq 1$. Note that under an orthogonal change of basis ($L_n \mapsto L'_n = U^T L_n U$, $|h\rangle \mapsto |h'\rangle = U^T |h\rangle$), Virasoro states transform into Virasoro states.

With reference to the vectors (2.12), the module \mathcal{V}_h is graded according to the level

$$\mathcal{V}_h = \bigoplus_{\ell=0}^{\infty} \mathcal{V}_{h,\ell}, \quad \ell = \sum_{i=1}^j n_i \quad (2.13)$$

The Virasoro character $\chi_h(q)$, which is the generating function for the spectrum of the Virasoro module \mathcal{V}_h , is

$$\chi_h(q) = \text{Tr}_{\mathcal{V}_h} q^{L_0 - c/24} = q^{-c/24+h} \sum_{\ell \in \mathbb{N}} d_\ell q^\ell, \quad d_\ell \geq 0 \quad (2.14)$$

where q is the modular parameter and the degeneracy $d_\ell = d_\ell^h = \dim \mathcal{V}_{h,\ell}$ is the dimension of the space of states at level ℓ .

2.3 Fields

Consider a bulk primary field $\phi^{(h)}(z, \bar{z})$. Radial quantization imposes the following definition of adjoint on the real surface $\bar{z} = z^*$

$$\phi^{(h)}(z, \bar{z})^\dagger = \bar{z}^{-2h} z^{-2\bar{h}} \phi^{(h)}(1/\bar{z}, 1/z) \quad (2.15)$$

1	$ 0\rangle$			
0	—			
q^2	$L_{-2} 0\rangle$			
q^3	$L_{-3} 0\rangle$			
$2q^4$	$L_{-4} 0\rangle$	$L_{-2}^2 0\rangle$		
$2q^5$	$L_{-5} 0\rangle$	$L_{-3}L_{-2} 0\rangle$		
$4q^6$	$L_{-6} 0\rangle$	$L_{-4}L_{-2} 0\rangle$	$L_{-3}^2 0\rangle$	$L_{-2}^3 0\rangle$
$4q^7$	$L_{-7} 0\rangle$	$L_{-5}L_{-2} 0\rangle$	$L_{-4}L_{-3} 0\rangle$	$L_{-3}L_{-2}^2 0\rangle$
$7q^8$	$L_{-8} 0\rangle$	$L_{-6}L_{-2} 0\rangle$	$L_{-5}L_{-3} 0\rangle$	$L_{-4}^2 0\rangle$
	$L_{-4}L_{-2}^2 0\rangle$	$L_{-3}^2L_{-2} 0\rangle$	$L_{-2}^4 0\rangle$	
$8q^9$
$12q^{10}$

Figure 1: Virasoro module \mathcal{V}_0 of Virasoro states in the vacuum $h = 0$ sector. The generic Virasoro character is $\chi_0(q) = \prod_{n=2}^{\infty} (1 - q^n)^{-1} = 1 + q^2 + q^3 + 2q^4 + 2q^5 + 4q^6 + 4q^7 + 7q^8 + 8q^9 + 12q^{10} + \dots$. In this sector, there is a null vector $L_{-1}|0\rangle = 0$ at level $\ell = 1$. For the minimal theories $\mathcal{M}(p', p)$, further null vectors appear. For example, for the Ising model $\mathcal{M}(3, 4)$, there is one null vector at level 6 and 7 and two at level 8. For $\mathcal{M}(4, 5)$, the first null vector enters at level 12. In practice it is often convenient to truncate the states at some fixed level.

for bulk (quasi-)primary fields (z and \bar{z} are independent complex variables, z^* is the complex conjugate of z). The decoupling between holomorphic and antiholomorphic degrees of freedom allows to drop the dependence on the antiholomorphic coordinate \bar{z} and consistently use chiral fields. The full picture is obtained after restoring the antiholomorphic dependence. The chiral form of the definition of adjoint

$$\phi^{(h)}(z)^\dagger = (z^*)^{-2h} \phi^{(h)}(1/z^*) \quad (2.16)$$

is used to define the dual space (bras)

$$\langle h| = \lim_{z \rightarrow 0} \langle 0| \phi^{(h)}(z)^\dagger = \lim_{u \rightarrow \infty} u^{2h} \langle 0| \phi^{(h)}(u). \quad (2.17)$$

Unlike the generators L_n , a chiral primary field $\phi^{(h)}(z)$ intertwines between different Virasoro modules, its action being dictated by the fusion rules

$$\phi^{(h_i)} \times \phi^{(h_j)} = \sum_k \mathcal{N}_{ij}^k \phi^{(h_k)} \quad (2.18)$$

If the fusion coefficient is $\mathcal{N}_{ij}^k \neq 0$ then the corresponding restriction of the field $\phi^{(h_i)}$, called chiral vertex operator (CVO), is non-zero

$$\phi \left(\begin{matrix} h_i \\ h_k, h_j \end{matrix} \right) (z) : \mathcal{V}_{h_j} \rightarrow \mathcal{V}_{h_k} \quad (2.19)$$

As a consequence of global conformal invariance we have the expression for the two-point function

$$\langle \phi^{(h)}(z) \phi^{(h)}(w) \rangle \equiv \langle 0| \phi^{(h)}(z) \phi^{(h)}(w) |0 \rangle = \frac{1}{(z - w)^{2h}} \quad (2.20)$$

corresponding to the trivial fusion coefficients $\mathcal{N}_{h_0}^h = \mathcal{N}_{hh}^0 = 1$ and to the non-zero CVOs:

$$\phi\left(\begin{matrix} h \\ h, 0 \end{matrix}\right)(z) : \quad \mathcal{V}_0 \rightarrow \mathcal{V}_h \quad (2.21)$$

$$\phi\left(\begin{matrix} h \\ 0, h \end{matrix}\right)(z) : \quad \mathcal{V}_h \rightarrow \mathcal{V}_0 \quad (2.22)$$

The nonvanishing of the generic fusion coefficient $\mathcal{N}_{hh_1}^{h_2} \neq 0$ corresponds to a non-zero three-point function

$$\begin{aligned} \langle \phi^{(h_2)}(u) \phi^{(h)}(z) \phi^{(h_1)}(w) \rangle &\equiv \langle 0 | \phi^{(h_2)}(u) \phi^{(h)}(z) \phi^{(h_1)}(w) | 0 \rangle \\ &= \frac{\sqrt{C_{h_2, h, h_1}}}{(u-z)^{h_2+h-h_1} (u-w)^{h_2+h_1-h} (z-w)^{h+h_1-h_2}} \end{aligned} \quad (2.23)$$

The two-point relation (2.20) is contained in (2.23) if we set h_1 or h_2 to zero with $C_{0, h, h_1} = C_{h_2, h, 0} = 1$.

Using (2.11) and (2.17) in the three-point function, we have

$$\langle h_2 | \phi^{(h)}(z) | h_1 \rangle = \lim_{\substack{u \rightarrow \infty \\ w \rightarrow 0}} u^{2h_2} \langle 0 | \phi^{(h_2)}(u) \phi^{(h)}(z) \phi^{(h_1)}(w) | 0 \rangle = \frac{\sqrt{C_{h_2, h, h_1}}}{z^{h+h_1-h_2}} \quad (2.24)$$

forcing the following Laurent expansion

$$\phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)(z) = z^{h_2-h_1-h} \sum_{k \in \mathbb{Z}} \phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)_{k+h_1-h_2} z^{-k} \quad (2.25)$$

with the operator modes

$$\phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)_{k+h_1-h_2} = \oint \frac{dz}{2\pi i} z^{k-1} \phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)(z) z^{h_1+h-h_2}, \quad k \in \mathbb{Z} \quad (2.26)$$

It is not possible to write a unique Laurent expansion for the various restrictions of a field. The analytic behaviour of $\phi^{(h)}(z) | h_1 \rangle$ as $z \rightarrow 0$ and of $\langle h_2 | \phi^{(h)}(z)$ as $z \rightarrow \infty$ requires the constraints

$$\begin{aligned} \phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)_{k+h_1-h_2} | h_1 \rangle &= 0, & k > 0 \\ \langle h_2 | \phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)_{k+h_1-h_2} &= 0, & k < 0 \end{aligned} \quad (2.27)$$

In particular, for the primary states the modes vanish if $k > 0$ giving

$$|h\rangle = \phi^{(h)}(0) | 0 \rangle = \phi\left(\begin{matrix} h \\ h, 0 \end{matrix}\right)_{-h} | 0 \rangle, \quad \langle h | h' \rangle = \delta_{h, h'} \quad (2.28)$$

Under infinitesimal coordinate transformations $z \mapsto z + \varepsilon_m z^{m+1}$, the primary fields must satisfy [2]

$$\delta_m \phi^{(h)}(z) = [L_m, \phi^{(h)}(z)] = z^{m+1} \frac{\partial}{\partial z} \phi^{(h)}(z) + h(m+1) z^m \phi^{(h)}(z) \quad (2.29)$$

and consequently the modes satisfy

$$[L_m, \phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)_{k+h_1-h_2}] = [-m(1-h) - (k+h_1-h_2)] \phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)_{m+k+h_1-h_2} \quad (2.30)$$

The chiral vertex operators (CVOs) are used in the quasi-particle approach [12] to construct quasi-particle states by repeated application of appropriate modes of primary fields

$$\phi\left(\begin{matrix} j_N \\ i_N, i_{N-1} \end{matrix}\right)_{-n_N-h_N+h_{N-1}} \cdots \phi\left(\begin{matrix} j_2 \\ i_2, i_1 \end{matrix}\right)_{-n_2-h_2+h_1} \phi\left(\begin{matrix} j_1 \\ i_1, 0 \end{matrix}\right)_{-n_1-h_1} |0\rangle \quad (2.31)$$

In general this set of states is overcomplete. It is supposed that a linearly independent set can be obtained by applying appropriate restrictions such as those arising from fusion and braiding. In this paper we concentrate on understanding the CVOs which are the building blocks of the quasi-particle approach.

2.4 Two and three-point functions

We can put the mode expansions inside the two-point function (2.20) and use the constraint (2.27) to reduce the number of terms

$$\langle 0 | \phi^{(h)}(z) \phi^{(h)}(w) | 0 \rangle = z^{-2h} \sum_{k \geq 0} \langle 0 | \phi\left(\begin{matrix} h \\ 0, h \end{matrix}\right)_{k+h} \phi\left(\begin{matrix} h \\ h, 0 \end{matrix}\right)_{-k-h} | 0 \rangle \left(\frac{w}{z}\right)^k \quad (2.32)$$

This expression is intimately related to the Laurent-Taylor expansion

$$(1-u)^{-2h} = \sum_{\ell \in \mathbb{Z}} a_\ell^h u^\ell, \quad a_\ell^h = 0 \text{ for } \ell < 0 \quad (2.33)$$

where

$$a_\ell^h = \langle 0 | \phi\left(\begin{matrix} h \\ 0, h \end{matrix}\right)_{\ell+h} \phi\left(\begin{matrix} h \\ h, 0 \end{matrix}\right)_{-\ell-h} | 0 \rangle \quad (2.34)$$

which we later verify for small $|\ell|$ from explicit matrices for the chiral modes. Of course the full result is obtained by restoring the antiholomorphic dependence

$$\langle \phi(z, \bar{z}) \phi(w, \bar{w}) \rangle = (z-w)^{-2h} (\bar{z}-\bar{w})^{-2h} \quad (2.35)$$

The structure constants and operator product expansions involve the calculation of three-point functions, that can be accessed in a completely similar way. The mode expansion in (2.24) together with (2.27) yields

$$\lim_{z \rightarrow 0} z^{h+h_1-h_2} \langle h_2 | \phi^{(h)}(z) | h_1 \rangle = \langle h_2 | \phi\left(\begin{matrix} h \\ h_2, h_1 \end{matrix}\right)_{h_1-h_2} | h_1 \rangle \equiv \sqrt{C_{h_2, h, h_1}} \quad (2.36)$$

where we see that the structure constant explicitly enters in the mode coupled with z^0 (fundamental mode).

The energy momentum tensor is not a primary field, but still its two-point function

$$\langle 0 | T(z) T(w) | 0 \rangle = \frac{c/2}{(z-w)^4} \quad (2.37)$$

is fixed by conformal invariance. We can proceed as we did for the fields. The mode expansion is now

$$\langle 0 | T(z) T(w) | 0 \rangle = z^{-4} \sum_{n \geq 0} \langle 0 | L_{n+2} L_{-n-2} | 0 \rangle \left(\frac{w}{z}\right)^n \quad (2.38)$$

The relevant Laurent-Taylor expansion is

$$(1-u)^{-4} = \sum_{\ell=0}^{\infty} a_\ell u^\ell = 1 + 4u + 10u^2 + 20u^3 + 35u^4 + 56u^5 + \cdots \quad (2.39)$$

where

$$\frac{c}{2} a_n = \langle 0 | L_{n+2} L_{-n-2} | 0 \rangle \quad (2.40)$$

Again we later verify this for small n from explicit matrices for the Virasoro generators.

3 Algorithm for Building Matrix Representations

In this section we present a general algorithm for building, level by level, matrix representations of the Virasoro generators and primary fields. We work in a distinguished basis which we call the L_1 -basis. The algorithm itself relates the vectors of the L_1 -basis to the basis of Virasoro states (2.12). Applying the algorithm for generic c also yields general expressions for the null vectors, level by level, for arbitrary minimal models $\mathcal{M}(p', p)$. Alternatively, these could be obtained level-by-level from the null vectors of the Gram matrix for Virasoro states.

We first present our general algorithm for the Virasoro generators in the L_1 -basis. In the subsequent subsection, we implement this algorithm for generic c in the $h = 0$ vacuum sector pointing out its salient features and discuss the relation between the L_1 -basis and the basis of Virasoro states. We implement our algorithm systematically for the critical Ising, tricritical Ising, 3-state Potts and Yang-Lee theories in Sections 4–7. In each case the algorithm was implemented using Mathematica [13].

3.1 Matrix algorithms for Virasoro generators in L_1 -basis

Because states belonging to different Virasoro modules are orthogonal, the matrices representing the Virasoro generators L_n can be obtained separately within each module \mathcal{V}_h

$$L_n = \oplus_h L_n^h \quad (3.1)$$

Usually we work within a fixed module and suppress the index h . To obtain matrix representations of Vir , we will take $L_{-n} = L_n^T$. For unitary theories L_n will be real so that $L_{-n} = L_n^T = L_n^\dagger$ whereas for non-unitary theories L_n will be complex with $L_{-n} = L_n^T \neq L_n^\dagger$. Suppose we can find a diagonal matrix L_0 and matrices $L_1 = L_{-1}^T$, $L_2 = L_{-2}^T$ satisfying

$$\begin{aligned} [L_1, L_0] &= L_1, & [L_2, L_0] &= 2L_2 \\ [L_1, L_{-1}] &= 2L_0, & [L_2, L_{-1}] &= 3L_1, & [L_2, L_{-2}] &= 4L_0 + c/2 \end{aligned} \quad (3.2)$$

Then we can extend this to a matrix representation of the full Virasoro algebra by recursively defining

$$L_n = L_{-n}^T = \frac{1}{n-2} [L_{n-1}, L_1], \quad n \geq 3 \quad (3.3)$$

Note that, if the generators L_n satisfy the Virasoro algebra, then the generators

$$L'_n = U^T L_n U, \quad U^T = U^{-1} \quad (3.4)$$

also satisfy the Virasoro algebra where $U = \oplus U^h$ is orthogonal and preserves the sectors (3.1). This corresponds to the freedom of an orthogonal change of basis. To specify a unique matrix representation of the Virasoro algebra one must first specify a distinguished *canonical* basis. The generators L_{-n} act naturally on the basis (2.12) of Virasoro states. In particular, the grading of these states imposes a block structure on the matrix representatives of the Virasoro generators

$$L_{-n}^h = \oplus_\ell L_{-n,\ell}^h, \quad L_{-n,\ell}^h : \mathcal{V}_{h,\ell} \rightarrow \mathcal{V}_{h,\ell+n} \quad (3.5)$$

The orthogonal matrix (3.4) must respect this grading so that

$$U^h = \bigoplus_{\ell=0}^{\infty} U_\ell^h \quad (3.6)$$

Within $\mathcal{V}_{h,\ell}$, the Virasoro states (2.12) with $\ell = \sum_{i=1}^j n_i$ are degenerate. However, they are not orthogonal and typically they contain null vectors, so a basis of Virasoro states (2.12) is not a good basis to choose in seeking matrix representations of the Virasoro generators. One possibility, is to use the Gram-Schmidt process to orthonormalize the Virasoro states with the null vectors removed. A better alternative, at least for unitary minimal models, is to use the basis of fermion states of FPI [4]. Since the fermion basis is a complete orthonormal basis of physical states we do not have to worry about removing null vectors. To obtain matrix representations of the Virasoro generators, we therefore start working in a fixed orthonormal basis of states.

With such a choice of basis, the structure of the matrices $L_1 = L_{-1}^T$ and $L_2 = L_{-2}^T$ is

$$L_1 = \begin{matrix} & d_0 & d_1 & d_2 & \cdots & d_{\ell-1} & & d_\ell \\ \begin{matrix} d_0 \\ d_1 \\ d_2 \\ \vdots \\ d_{\ell-1} \end{matrix} & \left(\begin{array}{cccccccc} \mathbf{0} & \star & \mathbf{0} & & & & & \\ & \mathbf{0} & \star & & & & & \\ & & \mathbf{0} & & & & & \\ & & & \ddots & & & & \\ & & & & \ddots & & & \\ & & & & & * & 0 & 0 & 0 & 0 & 0 & 0 \\ & & & & & 0 & * & 0 & 0 & 0 & 0 & 0 \\ & & & & & 0 & 0 & * & 0 & 0 & 0 & 0 \\ & & & & & 0 & 0 & 0 & * & 0 & 0 & 0 \end{array} \right) \end{matrix} \quad (3.7)$$

$$L_2 = \begin{matrix} & d_0 & d_1 & d_2 & \cdots & d_{\ell-1} & & d_\ell \\ \begin{matrix} d_0 \\ d_1 \\ d_2 \\ \vdots \\ d_{\ell-2} \end{matrix} & \left(\begin{array}{cccccccc} \mathbf{0} & \mathbf{0} & \star & & & & & \\ & \mathbf{0} & \mathbf{0} & \star & & & & \\ & & \mathbf{0} & \mathbf{0} & & & & \\ & & & \ddots & & & & \\ & & & & \ddots & & & \\ & & & & & * & * & * & * & * & 0 & 0 \\ & & & & & * & * & * & * & * & * & 0 \\ & & & & & * & * & * & * & * & * & * \end{array} \right) \end{matrix} \quad (3.8)$$

where $d_\ell = d_\ell^h = \dim \mathcal{V}_{h,\ell}$ as in (2.14), \star denotes non-zero off-diagonal blocks, the non-zero ℓ -blocks in position ℓ are of sizes $d_{\ell-1} \times d_\ell$ and $d_{\ell-2} \times d_\ell$ respectively and

$$\chi_h(q) = q^{-c/24+h} \sum_{\ell=0}^{\infty} d_\ell q^\ell \quad (3.9)$$

We allow for arbitrary entries in each non-zero block, substitute into the basic Virasoro commutators (3.2) and proceed to solve the equations level-by-level. At a given level, we find that the equations are underdetermined and only involve dot products between rows in the ℓ -blocks. This freedom corresponds to an allowed change of orthogonal basis implemented by the orthogonal matrix U_ℓ^h and is removed by demanding that the entries above the leading diagonal vanish in the $(d_{\ell-1} + d_{\ell-2}) \times d_\ell$ matrix B_ℓ formed by placing the ℓ block of L_1 above the ℓ -block of L_2 . In this basis, we find that the entries below the leading diagonal of the ℓ -block of L_1 automatically vanish. There remains a residual freedom associated with the choice of sign of each of the basis elements. We remove this by further demanding that the first non-zero entry at the top of each column of B_ℓ is positive. If this is pure imaginary, as can occur in non-unitary cases, we demand that the imaginary part of this non-zero entry of B_ℓ is positive. This algorithm thus fixes a distinguished *canonical* basis which we call the L_1 -basis because of the diagonal form of the blocks of L_1 . The L_1 -basis should be viewed as the fermionic basis of paths rotated by the orthogonal transformation U . Notice that only the level degeneracies, which can be read off from the Virasoro character, and the orthonormality of the starting basis is essential in this algorithm — the precise description of the starting basis is unimportant.

$$\begin{aligned}
&= \left(\sqrt{\frac{c}{2}} \right) \oplus \left(3\sqrt{\frac{2}{5}} \sqrt{c + \frac{22}{5}} \right) \oplus \left(\frac{7}{\sqrt{5}} \sqrt{\frac{1}{2}(c + \frac{22}{5})} \right) \oplus \left(\begin{array}{cccc} 8\sqrt{\frac{2}{7}} & \frac{1}{3}\sqrt{\frac{5}{2}(c + \frac{22}{5})} & \frac{1}{3}\sqrt{2(c + \frac{29}{70})} & 0 \\ 0 & 2 & \frac{4}{\sqrt{5}}\sqrt{\frac{c + \frac{22}{5}}{c + \frac{29}{70}}} & \sqrt{\frac{3(c - \frac{1}{2})(c + \frac{68}{7})}{2(c + \frac{29}{70})}} \end{array} \right) \quad (3.14) \\
&\oplus \left(\begin{array}{cccc} 9\sqrt{\frac{5}{14}} & \frac{1}{\sqrt{6}}\sqrt{c + \frac{22}{5}} & \frac{1}{\sqrt{3}}\sqrt{c + \frac{29}{70}} & 0 \\ 0 & \frac{13}{\sqrt{15}} & \frac{8}{\sqrt{30}}\sqrt{\frac{c + \frac{22}{5}}{c + \frac{29}{70}}} & \sqrt{\frac{(c - \frac{1}{2})(c + \frac{68}{7})}{(c + \frac{29}{70})}} \end{array} \right) \\
&\oplus \left(\begin{array}{cccccc} 5\sqrt{\frac{5}{3}} & \sqrt{\frac{7}{66}(c + \frac{22}{5})} & \sqrt{\frac{14}{39}(c + \frac{29}{70})} & 0 & \sqrt{\frac{5}{143}(c + \frac{11}{105})} & 0 \\ 0 & 14\sqrt{\frac{6}{55}} & 8\sqrt{\frac{3(c + \frac{22}{5})}{130(c + \frac{29}{70})}} & 3\sqrt{\frac{(c - \frac{1}{2})(c + \frac{68}{7})}{13(c + \frac{29}{70})}} & \frac{10}{3\sqrt{1001}}\sqrt{\frac{c + \frac{22}{5}}{c + \frac{11}{105}}} & \frac{1}{3}\sqrt{\frac{13(1050c^2 + 3305c - 251)}{105(c + \frac{11}{105})}} & 0 \\ 0 & 0 & 3\sqrt{\frac{6}{13}} & 0 & \frac{82}{3}\sqrt{\frac{11(c + \frac{29}{70})}{455(c + \frac{11}{105})}} & \frac{1}{30}\sqrt{\frac{13(c + \frac{22}{5})(1050c^2 + 3305c - 251)}{21(c + \frac{29}{70})(c + \frac{11}{105})}} & 0 \\ 0 & 0 & 0 & 3\sqrt{\frac{6}{13}} & 0 & 36\sqrt{\frac{70(c - \frac{1}{2})(c + \frac{68}{7})(c + \frac{11}{105})}{13(c + \frac{29}{70})(1050c^2 + 3305c - 251)}} & 10\sqrt{\frac{21(c + \frac{2}{5})(c + \frac{46}{3})(c + \frac{29}{70})}{(1050c^2 + 3305c - 251)}} \end{array} \right) \\
&\oplus \dots
\end{aligned}$$

We find that these matrices are generic in the sense that they can be reduced to produce the correct matrices for any special minimal value of c . More explicitly, if we specialize c to an allowed central charge for a minimal theory, say $c = 1/2$, then starting on the left we seek the first zero column of L_2 (excluding the first and third columns which are identically zero for all c). For $c = 1/2$ this occurs in column 11 corresponding to the known single null (unphysical) vector at level 6. We then remove the rows and columns of L_0 , L_1 and L_2 corresponding to this position and repeat the process. For $c = 1/2$, we find the next zero column of L_2 at the original position 15, corresponding to the single null vector at level 7, so we remove this row and column. At the next stage we remove rows and columns at the original positions 19 and 22, corresponding to the two null vectors at level 8, and so on. This process exactly produces the Virasoro generators of the critical Ising model $\mathcal{M}(3, 4)$ as given in Section 4. The same process works for the Yang-Lee theory with $c = -22/5$ as given in Section 7 and other special values of c corresponding to minimal models. Indeed, the matrix L_2 exhibits a remarkable structure of factors of the form $\sqrt{c - c_{p', p}}$ where $c_{p', p}$ is the central charge of the minimal model $\mathcal{M}(p', p)$.

The elements of the L_1 -basis can be obtained in terms of the Virasoro states from the generic minimal Virasoro matrices $L_{-n} = L_n^T$, $n = 1, 2$ by considering the action of these matrices on the fundamental basis formed by the column vectors $e_i = \{\delta_{i,j}\}_{j=1,2,\dots}$. We identify e_1 with the vacuum $|0\rangle$. Using the action of the Virasoro algebra on the states, the next states are identified as

$$L_{-2} e_1 \equiv L_{-2} |0\rangle = \sqrt{\frac{c}{2}} e_2 \Rightarrow e_2 = \sqrt{\frac{2}{c}} L_{-2} |0\rangle \quad (3.15)$$

$$L_{-1} e_2 \equiv \sqrt{\frac{2}{c}} L_{-3} |0\rangle = 2 e_3 \Rightarrow e_3 = \frac{1}{\sqrt{2c}} L_{-3} |0\rangle \quad (3.16)$$

This procedure continues for higher levels leading to the expressions in Figure 9. These expressions hold for both unitary and non-unitary models. The orthonormality of the L_1 basis is guaranteed by the orthonormality of the fundamental basis e_i .

3.3 Matrix algorithm for fields

An algorithm can similarly be applied level-by-level to obtain matrix representations of the primary fields $\phi^{(h)}(z)$. In the vector space $\mathcal{V}_{h_1} \oplus \mathcal{V}_{h_2}$ we consider the action of the two CVOs

$$\phi \left(\begin{array}{c} h \\ h_2, h_1 \end{array} \right) (z) : \quad \mathcal{V}_{h_1} \rightarrow \mathcal{V}_{h_2} \quad (3.17)$$

$$\phi \left(\begin{array}{c} h \\ h_1, h_2 \end{array} \right) (z) : \quad \mathcal{V}_{h_2} \rightarrow \mathcal{V}_{h_1} \quad (3.18)$$

$$\begin{array}{ll}
1 & |0\rangle \\
q^2 & \sqrt{\frac{2}{c}}L_{-2}|0\rangle \\
q^3 & \frac{1}{\sqrt{2c}}L_{-3}|0\rangle \\
2q^4 & \frac{1}{\sqrt{5c}}L_{-4}|0\rangle - \frac{\sqrt{2}(3L_{-4}-5L_{-2}^2)|0\rangle}{5\sqrt{c(c+\frac{22}{5})}} \\
2q^5 & \frac{1}{\sqrt{10c}}L_{-5}|0\rangle - \frac{(2L_{-5}-5L_{-3}L_{-2})|0\rangle}{5\sqrt{c(c+\frac{22}{5})}} \\
4q^6 & \sqrt{\frac{2}{35c}}L_{-6}|0\rangle - \frac{(8L_{-6}-10L_{-4}L_{-2}-5L_{-3}^2)|0\rangle}{15\sqrt{2c(c+\frac{22}{5})}} \\
4q^7 & \frac{(20L_{-6}+56L_{-4}L_{-2}-35L_{-3}^2)|0\rangle}{42\sqrt{10c(c+\frac{29}{70})}} - \frac{((78+60c)L_{-6}+(201+126c)L_{-4}L_{-2}-93L_{-3}^2-(29+70c)L_{-2}^3)|0\rangle}{35\sqrt{3c(c-\frac{1}{2})(c+\frac{22}{5})(c+\frac{68}{7})(c+\frac{29}{70})}} \\
& \dots \qquad \dots
\end{array}$$

Figure 2: L_1 -basis of orthonormal states for the generic Virasoro module \mathcal{V}_0 in terms of Virasoro states. The central charge c is treated as a parameter. Unphysical (null) states appear if c is set to the value of the central charge of a minimal model. For the Ising model with $c = 1/2$ an unphysical state appears at level $\ell = 6$. The apparent singularity of the state is an artifact of the normalization. Actually, the numerator is a null state. Further null states appear for the non-unitary cases with $c = -22/5$ and $c = -68/7$ corresponding to $\mathcal{M}(2, 5)$ and $\mathcal{M}(2, 7)$.

The Virasoro generators and fields have the following block structure when acting on $\mathcal{V}_{h_1} \oplus \mathcal{V}_{h_2}$

$$L_n = \begin{pmatrix} L_n^{h_1} & \mathbf{0} \\ \mathbf{0} & L_n^{h_2} \end{pmatrix}, \quad \phi^{(h)}(z) = \begin{pmatrix} \mathbf{0} & \phi\left(\begin{smallmatrix} h \\ h_1, h_2 \end{smallmatrix}\right)(z) \\ \phi\left(\begin{smallmatrix} h \\ h_2, h_1 \end{smallmatrix}\right)(z) & \mathbf{0} \end{pmatrix} \quad (3.19)$$

The two CVOs can be worked out separately because they do not mix under the commutation relation due to the block structure of the generators L_n . So we consider $\phi\left(\begin{smallmatrix} h \\ h_2, h_1 \end{smallmatrix}\right)(z)$ and expand the field in modes

$$\Phi(z) = \begin{pmatrix} \mathbf{0} & \mathbf{0} \\ \phi\left(\begin{smallmatrix} h \\ h_2, h_1 \end{smallmatrix}\right)(z) & \mathbf{0} \end{pmatrix}, \quad \Phi_k = \begin{pmatrix} \mathbf{0} & \mathbf{0} \\ \phi\left(\begin{smallmatrix} h \\ h_2, h_1 \end{smallmatrix}\right)_{k+h_1-h_2} & \mathbf{0} \end{pmatrix} \quad (3.20)$$

Working in the L_1 -basis for \mathcal{V}_{h_1} and \mathcal{V}_{h_2} , the block structure of the mode $\phi\left(\begin{smallmatrix} h \\ h_2, h_1 \end{smallmatrix}\right)_{h_1-h_2}$ is

$$\phi\left(\begin{smallmatrix} h \\ h_2, h_1 \end{smallmatrix}\right)_{h_1-h_2} = \begin{pmatrix} d_0^{h_2} \\ d_1^{h_2} \\ d_2^{h_2} \\ \vdots \\ d_\ell^{h_2} \end{pmatrix} \begin{pmatrix} d_0^{h_1} & d_1^{h_1} & d_2^{h_1} & \dots & d_\ell^{h_1} \\ \star & \mathbf{0} & \mathbf{0} & & \\ \mathbf{0} & \star & \mathbf{0} & & \\ \mathbf{0} & \mathbf{0} & \star & & \\ & & & \ddots & \\ & & & & * & * & * & * \\ & & & & * & * & * & * \\ & & & & * & * & * & * \\ & & & & * & * & * & * \\ & & & & * & * & * & * \\ & & & & & & & \ddots \end{pmatrix} \quad (3.21)$$

where

$$\chi_{h_1}(q) = q^{-c/24+h_1} \sum_{\ell=0}^{\infty} d_{\ell}^{h_1} q^{\ell}, \quad \chi_{h_2}(q) = q^{-c/24+h_2} \sum_{\ell=0}^{\infty} d_{\ell}^{h_2} q^{\ell}, \quad d_{\ell}^{h_1}, d_{\ell}^{h_2} \geq 0 \quad (3.22)$$

Given the mode Φ_0 we use (2.30) to generate the remaining modes by the action of the known matrix representation of L_n

$$\Phi_n = \frac{1}{n(h-1) + h_2 - h_1} [L_n, \Phi_0], \quad n \neq 0 \quad (3.23)$$

With the higher modes defined in this way, we solve the linear equations

$$[L_1, \Phi_1] = (h-2+h_2-h_1)\Phi_2, \quad [L_1, \Phi_{-1}] = (h+h_2-h_1)\Phi_0, \quad [L_2, \Phi_{-1}] = (2h-1+h_2-h_1)\Phi_1 \quad (3.24)$$

level-by-level for the entries of Φ_0 . We find that these three equations are consistent and sufficient to determine all the entries of the modes so that no higher level equations are required.

4 Critical Ising model

The critical Ising model $\mathcal{M}(3,4)$ with $c = 1/2$ is realized as the critical RSOS model on A_3 . It is essentially equivalent to the free fermion model and is well understood. We present it in this section to make contact with previous work and to illustrate the new features of our representation matrices in a familiar setting. In particular, we discuss the matrix representations of the Virasoro generators and fields in both the usual fermion basis and our L_1 -basis.

The free fermion (Ising) model is usually discussed in the context of a superconformal theory with Neveu-Schwarz ($h = 0, 1/2$) and Ramond ($h = 1/16$) sectors. The mode expansion of the free fermion field is given by

$$\psi(z) = \sum_k b_k z^{-k-1/2}, \quad k \in \begin{cases} \mathbb{Z} - 1/2, & \text{NS} \\ \mathbb{Z}, & \text{R} \end{cases} \quad (4.25)$$

with $b_{-k} = b_k^\dagger$. This is the primary field $\phi^{(1/2)}(z) = \psi(z)$ in the Neveu-Schwarz sector. The Virasoro generators have simple bilinear expressions in terms of normal ordered fermion modes

$$L_0 = \begin{cases} \sum_{k>0} k b_{-k} b_k, & \text{NS: } k \in \mathbb{Z} - \frac{1}{2} \\ \sum_{k>0} k b_{-k} b_k + \frac{1}{16}, & \text{R: } k \in \mathbb{Z} \end{cases} \quad (4.26)$$

$$L_n = \frac{1}{2} \sum_k (k + \frac{1}{2}) : b_{n-k} b_k : \quad : b_n b_m := \begin{cases} b_n b_m, & n < m \\ -b_m b_n, & n > m \end{cases} \quad (4.27)$$

with

$$[L_n, b_k] = (-\frac{n}{2} + k) b_{n+k} \quad (4.28)$$

Simple expressions for the generators L_n as bilinear forms only appear to hold for the Ising model.

$$T(z) = \begin{pmatrix} \frac{1}{16z^2} & \frac{1}{2\sqrt{2}z^3} & \frac{1}{\sqrt{2}z^4} & \frac{8}{\sqrt{51}z^5} & \frac{-7}{2\sqrt{102}z^5} & 8\sqrt{\frac{2}{51}}z^6 & \frac{-5}{\sqrt{51}z^6} & \frac{64}{\sqrt{935}z^7} & \frac{-20}{\sqrt{323}z^7} & \frac{-7}{2\sqrt{2090}z^7} & \frac{896}{\sqrt{115005}z^8} & -56\sqrt{\frac{10}{12397}}z^8 & \frac{5\sqrt{\frac{5}{1066}}}{9z^8} \\ \frac{1}{2\sqrt{2}z} & \frac{17}{16z^2} & \frac{3}{2z^3} & \frac{19}{\sqrt{102}z^4} & \frac{7}{\sqrt{51}z^4} & \frac{16}{\sqrt{51}z^5} & \frac{31}{2\sqrt{102}z^5} & \frac{56\sqrt{\frac{2}{935}}}{z^6} & \frac{25\sqrt{\frac{2}{323}}}{z^6} & \frac{21}{\sqrt{1045}z^6} & \frac{64\sqrt{\frac{22}{10455}}}{z^7} & \frac{116\sqrt{\frac{5}{12597}}}{z^7} & \frac{-40\sqrt{\frac{5}{533}}}{9z^7} \\ \frac{1}{\sqrt{2}} & \frac{3}{2z} & \frac{33}{16z^2} & \frac{\sqrt{\frac{51}{2}}}{2z^3} & 0 & \frac{7\sqrt{\frac{3}{17}}}{z^4} & \frac{\sqrt{\frac{6}{17}}}{z^4} & \frac{72\sqrt{\frac{2}{935}}}{z^5} & \frac{63}{2\sqrt{646}z^5} & \frac{-28}{\sqrt{1045}z^5} & \frac{296\sqrt{\frac{6}{38335}}}{z^6} & \frac{31\sqrt{\frac{15}{4199}}}{z^6} & \frac{80\sqrt{\frac{5}{533}}}{9z^6} \\ \frac{8}{\sqrt{51}} & \frac{19}{\sqrt{102}} & \frac{\sqrt{\frac{51}{2}}}{2z} & \frac{49}{16z^2} & 0 & \frac{5}{\sqrt{2}z^3} & 0 & \frac{17\sqrt{\frac{3}{55}}}{z^4} & \frac{\sqrt{\frac{3}{19}}}{z^4} & \frac{56\sqrt{\frac{2}{53295}}}{z^4} & \frac{208}{\sqrt{2255}z^5} & \frac{\sqrt{\frac{95}{26}}}{z^5} & \frac{-320\sqrt{\frac{10}{27183}}}{9z^5} \\ \frac{-7z}{2\sqrt{102}} & \frac{7}{\sqrt{51}} & 0 & 0 & \frac{49}{16z^2} & 0 & \frac{7}{2\sqrt{2}z^3} & 0 & \frac{7\sqrt{\frac{3}{38}}}{z^4} & \frac{13\sqrt{\frac{55}{969}}}{z^4} & 0 & \frac{56}{\sqrt{1235}z^5} & \frac{104\sqrt{\frac{110}{969}}}{9z^5} \\ 8\sqrt{\frac{2}{51}}z^2 & \frac{16z}{\sqrt{51}} & \frac{7\sqrt{\frac{3}{17}}}{\sqrt{17}} & \frac{5}{\sqrt{2}z} & 0 & \frac{65}{16z^2} & 0 & \frac{\sqrt{\frac{165}{2}}}{2z^3} & 0 & 0 & \frac{67\sqrt{\frac{5}{902}}}{z^4} & \frac{2\sqrt{\frac{5}{247}}}{z^4} & \frac{128\sqrt{\frac{5}{27183}}}{9z^4} \\ \frac{-5z^2}{\sqrt{51}} & \frac{31z}{2\sqrt{102}} & \frac{7\sqrt{\frac{2}{935}}}{\sqrt{935}} & 17\sqrt{\frac{3}{55}}z & \frac{7}{2\sqrt{2}z} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{-31\sqrt{\frac{82}{3315}}}{9z^4} \\ \frac{64z}{\sqrt{935}} & \frac{16z}{\sqrt{51}} & \frac{7\sqrt{\frac{2}{935}}}{\sqrt{935}} & \frac{63z}{2\sqrt{646}} & \frac{7\sqrt{\frac{3}{38}}}{\sqrt{38}} & 0 & \frac{\sqrt{57}}{2z} & 0 & \frac{\sqrt{57}}{2z^3} & 0 & 0 & \frac{\sqrt{\frac{105}{2}}}{2z^3} & 0 \\ \frac{-20z^3}{\sqrt{323}} & \frac{21z^2}{\sqrt{1045}} & \frac{-28z}{\sqrt{1045}} & \frac{56\sqrt{\frac{2}{53295}}}{\sqrt{2255}} & \frac{13\sqrt{\frac{55}{969}}}{\sqrt{969}} & 0 & 0 & 0 & 0 & \frac{81}{16z^2} & 0 & 0 & 0 \\ \frac{-7z^3}{2\sqrt{2090}} & \frac{21z^2}{\sqrt{1045}} & \frac{-28z}{\sqrt{1045}} & \frac{208z}{\sqrt{2255}} & 0 & 0 & 0 & 0 & 0 & \frac{81}{16z^2} & 0 & 0 & 0 \\ \frac{896z^4}{\sqrt{115005}} & \frac{64\sqrt{\frac{22}{10455}}z^3}{\sqrt{10455}} & \frac{296\sqrt{\frac{6}{38335}}z^2}{\sqrt{38335}} & \frac{208z}{\sqrt{2255}} & 0 & \frac{67\sqrt{\frac{5}{902}}}{z^4} & 0 & \frac{\sqrt{123}}{2z} & 0 & 0 & \frac{97}{16z^2} & 0 & 0 \\ -56z^4\sqrt{\frac{10}{12597}} & \frac{116\sqrt{\frac{5}{12597}}z^3}{\sqrt{12597}} & \frac{31\sqrt{\frac{15}{4199}}z^2}{\sqrt{4199}} & \frac{208z}{\sqrt{2255}} & \frac{56z}{\sqrt{1235}} & \frac{2\sqrt{\frac{5}{247}}}{\sqrt{2470}} & \frac{163}{\sqrt{2470}} & 0 & \frac{\sqrt{195}}{2z} & 0 & 0 & \frac{97}{16z^2} & 0 \\ \frac{-49z^4}{9\sqrt{1045}} & \frac{91\sqrt{\frac{11}{190}}z^3}{\sqrt{190}} & \frac{-217\sqrt{\frac{2}{1045}}z^2}{\sqrt{1045}} & \frac{-28\sqrt{\frac{19}{2805}}z^2}{\sqrt{2805}} & \frac{104\sqrt{\frac{110}{969}}z}{\sqrt{969}} & \frac{112\sqrt{\frac{10}{10659}}}{9} & \frac{91\sqrt{\frac{55}{969}}}{9} & 0 & 0 & \frac{9}{2\sqrt{2}z^3} & 0 & \frac{97}{16z^2} & 0 \\ \frac{5\sqrt{\frac{5}{1066}}z^4}{9} & \frac{-40\sqrt{\frac{5}{533}}z^3}{9} & \frac{80\sqrt{\frac{5}{533}}z^2}{9} & \frac{-320\sqrt{\frac{10}{27183}}z}{9} & \frac{217\sqrt{\frac{41}{3315}}z}{18} & \frac{128\sqrt{\frac{5}{3315}}}{9} & \frac{-31\sqrt{\frac{82}{3315}}}{9} & 0 & 0 & 0 & 0 & 0 & \frac{97}{16z^2} \end{pmatrix}$$

Figure 3: Energy-momentum tensor (in the L_1 basis) of the critical Ising model in the $h = 1/16$ sector.

$$\phi\left(\begin{matrix} 1/2 \\ 1/2, 0 \end{matrix}\right)(z) = \begin{pmatrix} 1 & z^{-2} & z^{-3} & \frac{3}{\sqrt{10}z^4} & \frac{1}{\sqrt{10}z^4} & \frac{2}{\sqrt{5}z^5} & \frac{1}{\sqrt{5}z^5} & \frac{\sqrt{5}}{z^6} & \frac{1}{3\sqrt{14}z^6} & \frac{3}{\sqrt{14}z^7} & \frac{1}{\sqrt{3}z^7} & \frac{1}{\sqrt{42}z^7} \\ z & -\frac{1}{z} & 0 & \frac{1}{\sqrt{10}z^3} & \frac{-3}{\sqrt{10}z^3} & \frac{-2}{\sqrt{5}z^4} & \frac{-2}{\sqrt{5}z^4} & \frac{3}{\sqrt{35}z^5} & \frac{-5}{3\sqrt{14}z^5} & \frac{\sqrt{7}}{z^6} & -\frac{1}{\sqrt{3}z^6} & -\frac{2}{z^6}\sqrt{\frac{2}{21}} \\ z^2 & 0 & -\frac{1}{z} & -\frac{1}{\sqrt{10}z^2} & \frac{3}{\sqrt{10}z^2} & 0 & 0 & \frac{1}{\sqrt{35}z^4} & \frac{5\sqrt{2}}{3z^4} & \frac{1}{\sqrt{14}z^5} & -\frac{1}{\sqrt{3}z^5} & \frac{5}{\sqrt{42}z^5} \\ z^3 & 0 & 0 & \frac{-3}{\sqrt{10}z} & -\frac{1}{\sqrt{10}z} & \frac{2}{\sqrt{5}z^2} & \frac{2}{\sqrt{5}z^2} & -\frac{1}{\sqrt{35}z^3} & \frac{-5\sqrt{2}}{3z^3} & 0 & 0 & 0 \\ z^4 & 0 & 0 & 0 & 0 & \frac{-2}{\sqrt{5}z} & \frac{-2}{\sqrt{5}z} & \frac{-3}{\sqrt{35}z^2} & \frac{7}{3\sqrt{10}z^2} & -\frac{1}{\sqrt{14}z^3} & \frac{1}{\sqrt{3}z^3} & \frac{-5}{\sqrt{42}z^3} \\ 0 & z^2 & -z & \frac{1}{\sqrt{10}} & \frac{-3}{\sqrt{10}} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ z^5 & 0 & 0 & 0 & 0 & 0 & 0 & -\frac{\sqrt{5}}{z} & \frac{-1}{3\sqrt{14}z} & -\frac{\sqrt{2}}{z^2} & \frac{1}{\sqrt{3}z^2} & \frac{2}{z^2}\sqrt{\frac{2}{21}} \\ 0 & z^3 & 0 & \frac{-3z}{\sqrt{10}} & -\frac{z}{\sqrt{10}} & \frac{-2}{\sqrt{5}} & \frac{-2}{\sqrt{5}} & 0 & 0 & 0 & 0 & 0 \\ z^6 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{-3}{\sqrt{14}z} & -\frac{1}{\sqrt{3}z} & -\frac{1}{\sqrt{42}z} \\ 0 & \frac{2z^4}{\sqrt{5}} & \frac{z^3}{\sqrt{5}} & \frac{-3z^2}{5\sqrt{2}} & \frac{-z^2}{5\sqrt{2}} & \frac{-2z}{5} & \frac{-2z}{5} & \frac{\sqrt{7}}{5} & 0 & 0 & 0 & 0 \\ 0 & \frac{z^4}{\sqrt{5}} & \frac{-2z^3}{\sqrt{5}} & \frac{3\sqrt{3}z^2}{5} & \frac{\sqrt{2}z^2}{5} & \frac{-2z}{5} & \frac{-2z}{5} & \frac{1}{5\sqrt{7}} & \frac{-5\sqrt{5}}{3} & 0 & 0 & 0 \\ z^7 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2\sqrt{\frac{5}{33}}z^5 & \frac{8z^4}{\sqrt{165}} & \frac{z^3}{5\sqrt{66}} & -\frac{\sqrt{3}z^3}{5} & \frac{-17z^2}{5\sqrt{33}} & \frac{-2\sqrt{\frac{3}{11}}z^2}{5} & -\frac{7}{5}\sqrt{\frac{7}{33}}z & 0 & 2\sqrt{\frac{14}{165}} & \frac{-6}{\sqrt{55}} & 0 \\ 0 & \sqrt{\frac{5}{13}}z^5 & \frac{-6z^4}{\sqrt{65}} & -\frac{\sqrt{2}z^3}{5} & \frac{3\sqrt{\frac{2}{13}}z^3}{5} & \frac{14z^2}{5\sqrt{13}} & \frac{2z^2}{5\sqrt{13}} & \frac{-27z}{5\sqrt{91}} & \frac{-5\sqrt{\frac{5}{182}}z}{3} & 2\sqrt{\frac{2}{455}} & \frac{1}{\sqrt{195}} & -5\sqrt{\frac{10}{273}} \\ 0 & \frac{2z^5}{\sqrt{429}} & \frac{-5z^4}{\sqrt{429}} & 2\sqrt{\frac{10}{429}}z^3 & -2\sqrt{\frac{30}{143}}z^3 & -2\sqrt{\frac{5}{429}}z^2 & 3\sqrt{\frac{15}{143}}z^2 & 2\sqrt{\frac{5}{3003}}z & \sqrt{\frac{66}{91}}z & -\frac{1}{\sqrt{6006}} & \frac{1}{\sqrt{143}} & -\sqrt{\frac{11}{182}} \end{pmatrix} \dots$$

Figure 4: The lower-left block $\phi\left(\begin{matrix} 1/2 \\ 1/2, 0 \end{matrix}\right)(z)$ of the primary field $\phi^{(1/2)}(z)$ (in the L_1 basis) of the critical Ising model. Notice that the squares of the coefficients in the first column give the expansion of $(1-u)^{-1}$ in agreement with the two-point function (2.33).

$$\phi\left(\begin{matrix} 1/16 \\ 1/16, 0 \end{matrix}\right)(z) = \left(\begin{array}{cccccccc} 1 & \frac{1}{8z^2} & \frac{1}{8z^3} & \frac{3}{8\sqrt{10}z^4} & \frac{1}{4\sqrt{5}z^5} & \frac{3}{64\sqrt{5}z^5} & \frac{1}{8z^6} & \frac{5}{64z^6} \\ z & \frac{-15}{16\sqrt{2}z} & \frac{-7}{16\sqrt{2}z^2} & \frac{-13}{32\sqrt{5}z^3} & \frac{-3}{8\sqrt{10}z^4} & \frac{-69}{128\sqrt{10}z^4} & \frac{-11}{16\sqrt{70}z^5} & \frac{\sqrt{5}}{64z^5} \\ 3z^2 & \frac{35}{64\sqrt{2}} & \frac{-45}{64\sqrt{2}z} & \frac{-87}{128\sqrt{5}z^2} & \frac{-21}{32\sqrt{10}z^3} & \frac{217}{512\sqrt{10}z^3} & \frac{-81}{64\sqrt{70}z^4} & \frac{3055}{8192\sqrt{7}z^4} \\ 8\sqrt{2} & \frac{35z^2}{64\sqrt{2}} & \frac{-45z}{64\sqrt{2}z} & \frac{-9\sqrt{255}}{256z} & \frac{-11\sqrt{51}}{128z^2} & \frac{2048z^2}{2048z^2} & \frac{-43\sqrt{51}}{256z^3} & \frac{79}{1024\sqrt{5}z^4} \\ \frac{\sqrt{51}z^3}{32} & \frac{35z}{256\sqrt{51}} & \frac{875}{256\sqrt{51}} & \frac{256z}{256z} & \frac{128z^2}{128z^2} & \frac{2048z^2}{2048z^2} & \frac{256z^3}{256z^3} & \frac{-94705}{16384\sqrt{714}z^3} \\ 0 & \frac{7z}{\sqrt{102}} & \frac{-7}{2\sqrt{102}} & 0 & 0 & \frac{-5\sqrt{51}}{16z^2} & \frac{-5\sqrt{51}}{16z^2} & \frac{-49\sqrt{7}}{128z^3} \\ 5\sqrt{\frac{51}{2}}z^4 & \frac{21\sqrt{\frac{3}{34}}z^2}{1024} & \frac{287z}{1024\sqrt{102}} & \frac{17213}{2048\sqrt{255}} & \frac{-15\sqrt{255}}{512z} & \frac{-3565\sqrt{\frac{15}{34}}}{8192z} & \frac{4991\sqrt{\frac{5}{51}}}{16384z^2} & \frac{435643}{131072\sqrt{357}z^2} \\ 0 & \frac{3\sqrt{\frac{3}{17}}z^2}{4} & \frac{31z}{8\sqrt{51}} & \frac{-\sqrt{10}}{51} & 0 & \frac{-35\sqrt{\frac{15}{17}}}{64z} & \frac{-35\sqrt{\frac{5}{102}}}{32z^2} & \frac{161\sqrt{\frac{7}{102}}}{256z^2} \\ \frac{3\sqrt{935}z^5}{512} & \frac{21\sqrt{\frac{17}{56}}z^3}{4096} & \frac{1197z^2}{4096\sqrt{935}} & \frac{13391z}{20480\sqrt{374}} & \frac{7777}{10240\sqrt{187}} & \frac{204631}{163840\sqrt{187}} & \frac{-225\sqrt{\frac{187}{7}}}{4096z} & \frac{-435643\sqrt{\frac{5}{2018}}}{262144z} \\ 0 & \frac{3\sqrt{\frac{17}{19}}z^3}{16} & \frac{189z^2}{32\sqrt{323}} & \frac{5\sqrt{5}}{646z} & \frac{-4\sqrt{5}}{323} & \frac{565\sqrt{\frac{5}{323}}}{256} & \frac{-35\sqrt{\frac{95}{34}}}{128z} & \frac{-805\sqrt{\frac{7}{646}}}{1024z} \\ 0 & \frac{7z^3}{\sqrt{2090}} & \frac{-21z^2}{2\sqrt{2090}} & \frac{21z}{10\sqrt{209}} & \frac{-7}{10\sqrt{418}} & \frac{-373}{80\sqrt{418}} & 0 & \frac{-189\sqrt{\frac{35}{209}}}{128z} \\ \frac{\sqrt{115005}z^6}{2048} & \frac{35\sqrt{\frac{255}{451}}z^4}{16384} & \frac{511\sqrt{\frac{51}{2255}}z^3}{16384} & \frac{72093\sqrt{\frac{3}{15334}}z^2}{81920} & \frac{21679\sqrt{\frac{11}{2091}}z}{40960} & \frac{94829\sqrt{\frac{3}{7667}}z}{655360} & \frac{2699239\sqrt{\frac{7}{23001}}}{81920} & \frac{16684913}{655360\sqrt{46002}} \\ 0 & \frac{5\sqrt{\frac{255}{494}}z^4}{32} & \frac{\sqrt{\frac{4845}{26}}z^3}{64} & \frac{93\sqrt{\frac{3}{4199}}z^2}{8} & \frac{29\sqrt{\frac{2}{12597}}z}{9} & \frac{295\sqrt{\frac{3}{8398}}z}{512} & \frac{-16\sqrt{\frac{14}{12597}}}{9} & \frac{1127\sqrt{\frac{35}{12597}}}{4096} \\ 0 & \frac{35\sqrt{\frac{5}{209}}z^4}{36} & \frac{-7\sqrt{\frac{10}{53}}z^3}{72} & \frac{-217z^2}{60\sqrt{418}} & \frac{91\sqrt{\frac{11}{19}}z}{360} & \frac{2611z}{320\sqrt{209}} & \frac{-14\sqrt{\frac{7}{209}}}{45} & \frac{147\sqrt{\frac{35}{418}}}{256} \\ 0 & \frac{5\sqrt{\frac{10}{533}}z^4}{9} & \frac{-10\sqrt{\frac{10}{533}}z^3}{9} & \frac{10z^2}{3\sqrt{533}} & \frac{-10\sqrt{\frac{2}{533}}z}{9} & \frac{-31z}{4\sqrt{1066}} & \frac{5\sqrt{\frac{2}{3731}}}{9} & \frac{279\sqrt{\frac{5}{3731}}}{64} \end{array} \right)$$

Figure 6: The left-central block $\phi\left(\begin{matrix} 1/16 \\ 1/16, 0 \end{matrix}\right)(z)$ of the primary field $\phi^{(1/16)}(z)$ (in the L_1 basis) of the critical Ising model. Notice that the squares of the coefficients in the first column give the expansion of $(1-u)^{-1/8}$ in agreement with the two-point function (2.33).

$$\phi_{\left(\frac{1}{2}, 1/16\right)}(z) = z^{3/8} \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{-3}{4z} & \frac{-5}{16z^2} & \frac{-65}{32\sqrt{102}z^3} & \frac{1}{-\sqrt{51}z^3} & \frac{-91\sqrt{\frac{3}{17}}}{256z^4} & \frac{-\sqrt{\frac{3}{34}}}{2z^4} & \frac{-2639}{512\sqrt{1870}z^5} & \frac{-29}{8\sqrt{646}z^5} & \frac{1}{-\sqrt{1045}z^5} \\ z & \frac{5}{8} & \frac{-21}{32z} & \frac{-185}{64\sqrt{102}z^2} & \frac{5}{2\sqrt{51}z^2} & \frac{-689}{512\sqrt{51}z^3} & \frac{7}{4\sqrt{102}z^3} & \frac{-6279}{1024\sqrt{1870}z^4} & \frac{33}{16\sqrt{646}z^4} & \frac{9}{2\sqrt{1045}z^4} \\ \frac{3z^2}{8\sqrt{2}} & \frac{7z}{32} & \frac{81}{128} & \frac{-529\sqrt{\frac{3}{34}}}{256z} & \frac{-5\sqrt{\frac{3}{17}}}{8z} & \frac{-3955}{2048\sqrt{51}z^2} & \frac{35}{16\sqrt{102}z^2} & \frac{-31837}{4096\sqrt{1870}z^3} & \frac{301}{64\sqrt{646}z^3} & \frac{-63}{8\sqrt{1045}z^3} \\ \frac{5z^3}{16\sqrt{2}} & \frac{9z^2}{64} & \frac{55z}{256} & \frac{3283}{512\sqrt{102}} & \frac{5}{16\sqrt{51}} & \frac{-5735\sqrt{\frac{3}{17}}}{4096z} & \frac{-35\sqrt{\frac{3}{34}}}{32z} & \frac{-18391\sqrt{\frac{5}{374}}}{8192z^2} & \frac{385}{128\sqrt{646}z^2} & \frac{21\sqrt{\frac{5}{209}}}{16z^2} \\ \frac{35z^4}{128\sqrt{2}} & \frac{55z^3}{512} & \frac{273z^2}{2048} & \frac{8893z}{4096\sqrt{102}} & \frac{5z}{128\sqrt{51}} & \frac{148525}{32768\sqrt{51}} & \frac{175}{256\sqrt{102}} & \frac{-323337\sqrt{\frac{5}{374}}}{65536z} & \frac{-9135}{1024\sqrt{646}z} & \frac{-63\sqrt{\frac{5}{209}}}{128z} \\ \frac{z^4}{128\sqrt{2}} & \frac{-35z^3}{512} & \frac{315z^2}{2048} & \frac{-1995\sqrt{\frac{3}{34}}z}{4096} & \frac{189\sqrt{\frac{3}{17}}z}{128} & \frac{4389\sqrt{\frac{3}{17}}}{-32768} & \frac{-297\sqrt{\frac{3}{34}}}{256} & \frac{-1197\sqrt{\frac{11}{170}}}{-65536z} & \frac{891}{1024\sqrt{646}z} & \frac{-123\sqrt{\frac{11}{95}}}{128z} \\ \frac{63z^5}{256\sqrt{2}} & \frac{91z^4}{1024} & \frac{405z^3}{4096} & \frac{3635\sqrt{\frac{3}{34}}z^2}{8192} & \frac{\sqrt{\frac{3}{17}}z^2}{256} & \frac{100849z}{65536\sqrt{51}} & \frac{49z}{512\sqrt{102}} & \frac{3590887}{131072\sqrt{1870}} & \frac{4487}{2048\sqrt{646}} & \frac{63}{256\sqrt{1045}} \\ \frac{3z^5}{256\sqrt{2}} & \frac{-81z^4}{1024} & \frac{385z^3}{4096} & \frac{1575\sqrt{\frac{3}{34}}z^2}{8192} & \frac{189\sqrt{\frac{3}{17}}z^2}{256} & \frac{-25137\sqrt{\frac{3}{17}}z}{65536} & \frac{783\sqrt{\frac{3}{34}}z}{512} & \frac{44289\sqrt{\frac{11}{170}}}{-131072} & \frac{-22869}{2048\sqrt{646}} & \frac{41\sqrt{\frac{1}{95}}}{256} \end{pmatrix}.$$

Figure 7: The bottom-central block $\phi_{\left(\frac{1}{2}, 1/16\right)}(z)$ of the primary field $\phi^{(1/16)}(z)$ (in the L_1 basis) of the critical Ising model.

$$\begin{aligned}
L_1 = & \left(\frac{1}{\sqrt{5}}\right) \oplus \left(2\sqrt{\frac{3}{5}}\right) \oplus \left(\sqrt{\frac{33}{5}} \ 0\right) \oplus \begin{pmatrix} \frac{8}{\sqrt{5}} & 0 & 0 \\ 0 & \sqrt{\frac{31}{5}} & 0 \end{pmatrix} \oplus \begin{pmatrix} \sqrt{21} & 0 & 0 & 0 \\ 0 & 6\sqrt{\frac{2}{5}} & 0 & 0 \\ 0 & 0 & \sqrt{\frac{41}{5}} & 0 \end{pmatrix} \oplus \\
& \begin{pmatrix} 2\sqrt{\frac{39}{5}} & 0 & 0 & 0 & 0 & 0 \\ 0 & \sqrt{\frac{123}{5}} & 0 & 0 & 0 & 0 \\ 0 & 0 & 2\sqrt{\frac{23}{5}} & 0 & 0 & 0 \\ 0 & 0 & 0 & \sqrt{\frac{51}{5}} & 0 & 0 \end{pmatrix} \oplus \begin{pmatrix} \sqrt{\frac{217}{5}} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2\sqrt{\frac{46}{5}} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 3\sqrt{\frac{17}{5}} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 4\sqrt{\frac{7}{5}} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \sqrt{\frac{61}{5}} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \sqrt{\frac{61}{5}} & 0 & 0 \end{pmatrix} \\
& \oplus \begin{pmatrix} 12\sqrt{\frac{2}{5}} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \sqrt{51} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 4\sqrt{\frac{14}{5}} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \sqrt{\frac{183}{5}} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 2\sqrt{\frac{33}{5}} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 2\sqrt{\frac{33}{5}} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & \sqrt{\frac{71}{5}} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \sqrt{\frac{71}{5}} & 0 & 0 & 0 \end{pmatrix} \oplus \dots \tag{5.63}
\end{aligned}$$

$$L_2 = \begin{pmatrix} \dots & \frac{\sqrt{3}}{2} & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \frac{13}{2\sqrt{11}} & \sqrt{\frac{10}{11}} & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \frac{23}{4}\sqrt{\frac{3}{11}} & 2\sqrt{\frac{30}{341}} & \frac{1}{4}\sqrt{\frac{65}{31}} & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \frac{11}{4}\sqrt{\frac{15}{7}} & \sqrt{\frac{5}{31}} & \frac{1}{4}\sqrt{\frac{2145}{1271}} & 5\sqrt{\frac{15}{3157}} \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & 0 & \frac{1}{2}\sqrt{\frac{31}{2}} & 0 & \frac{1}{2}\sqrt{\frac{861}{22}} \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \end{pmatrix} \tag{5.64}$$

$$\begin{aligned}
L_2 = & \left(\frac{\sqrt{3}}{2}\right) \oplus \left(\frac{13}{2\sqrt{11}} \ \sqrt{\frac{10}{11}}\right) \oplus \left(\frac{23}{4}\sqrt{\frac{3}{11}} \ 2\sqrt{\frac{30}{341}} \ \frac{1}{4}\sqrt{\frac{65}{31}}\right) \oplus \begin{pmatrix} \frac{11}{4}\sqrt{\frac{15}{7}} & \sqrt{\frac{5}{31}} & \frac{1}{4}\sqrt{\frac{2145}{1271}} & 5\sqrt{\frac{15}{3157}} \\ 0 & \frac{1}{2}\sqrt{\frac{31}{2}} & 0 & \frac{1}{2}\sqrt{\frac{861}{22}} \end{pmatrix} \oplus \\
& \begin{pmatrix} \frac{43}{2}\sqrt{\frac{5}{91}} & 8\sqrt{\frac{5}{3813}} & \sqrt{\frac{2145}{29233}} & 40\sqrt{\frac{5}{53669}} & \frac{5}{4}\sqrt{\frac{274355}{6877299}} & 0 \\ 0 & \frac{103}{2\sqrt{246}} & 0 & \frac{1}{2}\sqrt{\frac{8897}{374}} & -3610\sqrt{\frac{598}{39124724001}} & 3\sqrt{\frac{32890}{1701001}} \\ 0 & 0 & \frac{3}{2}\sqrt{\frac{41}{23}} & 0 & \frac{2067}{4}\sqrt{\frac{697}{39123023}} & 200\sqrt{\frac{345}{1701001}} \end{pmatrix} \oplus \tag{5.65} \\
& \begin{pmatrix} \frac{53}{2}\sqrt{\frac{21}{403}} & 10\sqrt{\frac{14}{29233}} & 5\sqrt{\frac{1001}{496961}} & 50\sqrt{\frac{3}{53669}} & \frac{25}{4}\sqrt{\frac{384097}{139838413}} & 0 & \frac{125}{2}\sqrt{\frac{239}{56932337}} & 0 \\ 0 & 113\sqrt{\frac{3}{1886}} & 0 & \frac{3}{4}\sqrt{\frac{1271}{187}} & -14440\sqrt{\frac{897}{795536054687}} & 36\sqrt{\frac{16445}{103761061}} & \frac{245}{2}\sqrt{\frac{51}{3688487}} & \frac{13}{4}\sqrt{\frac{165}{451949}} \\ 0 & 0 & \frac{133}{2\sqrt{391}} & 0 & \frac{84747}{4}\sqrt{\frac{17}{2386504403}} & 200\sqrt{\frac{14145}{103761061}} & \frac{-95}{2}\sqrt{\frac{299}{10161563}} & 486\sqrt{\frac{65}{18529909}} \\ 0 & 0 & 0 & \frac{3}{4}\sqrt{\frac{51}{7}} & 0 & 0 & \frac{3}{2}\sqrt{\frac{739381}{68593}} & \frac{13}{4}\sqrt{\frac{2135}{9799}} \end{pmatrix}
\end{aligned}$$

The energy momentum tensor is shown in Figure 9.

Taking the direct sum of the two previous Virasoro modules, we can extract the field $\phi^{(1/10)}(z)$ whose block structure is shown in (3.19). We present here the lower-left 10×10 block of the

field:

$$\phi\left(\frac{1}{10}, 0\right)(z) = \quad (5.66)$$

$$\left(\begin{array}{cccccccccccc} 1 & \frac{1}{\sqrt{35}z^2} & \frac{1}{\sqrt{35}z^3} & \frac{3}{5\sqrt{14}z^4} & \frac{1}{5z^4}\sqrt{\frac{3}{238}} & \frac{2}{5\sqrt{7}z^5} & \frac{1}{5z^5}\sqrt{\frac{3}{119}} & \frac{1}{7z^6} & \frac{1}{\sqrt{714}z^6} & \frac{2}{35\sqrt{195}z^6} \\ \frac{z}{\sqrt{5}} & \frac{-9}{5\sqrt{7}z} & \frac{-4}{5\sqrt{7}z^2} & \frac{-1}{5z^3}\sqrt{\frac{7}{10}} & \frac{-19}{5z^3}\sqrt{\frac{3}{1190}} & \frac{-3}{5\sqrt{35}z^4} & \frac{-2}{5z^4}\sqrt{\frac{21}{85}} & \frac{-1}{7\sqrt{5}z^5} & \frac{-11}{\sqrt{3570}z^5} & \frac{-58}{175\sqrt{39}z^5} \\ \frac{\sqrt{3}z^2}{5} & \frac{6}{5}\sqrt{\frac{3}{35}} & \frac{-9}{5z}\sqrt{\frac{3}{35}} & \frac{-17}{25z^2}\sqrt{\frac{3}{14}} & \frac{19}{25z^2}\sqrt{\frac{7}{34}} & \frac{-8}{25z^3}\sqrt{\frac{3}{7}} & \frac{38}{25\sqrt{119}z^3} & \frac{-3\sqrt{3}}{35z^4} & \frac{1}{5z^4}\sqrt{\frac{2}{119}} & \frac{232}{175\sqrt{65}z^4} \\ \frac{z^3}{5}\sqrt{\frac{11}{5}} & \frac{6z}{25\sqrt{77}} & \frac{96}{25\sqrt{77}} & \frac{-27}{25z}\sqrt{\frac{11}{70}} & \frac{-57}{25z^2}\sqrt{\frac{21}{1870}} & \frac{-13}{25z^2}\sqrt{\frac{11}{35}} & \frac{38}{25z^2}\sqrt{\frac{21}{935}} & \frac{-1}{7z^3}\sqrt{\frac{11}{5}} & \frac{19}{z^3}\sqrt{\frac{2}{19635}} & \frac{-4408}{875\sqrt{429}z^3} \\ 0 & 2\sqrt{\frac{10}{77}z} & -\sqrt{\frac{10}{77}} & 0 & \frac{-12}{z}\sqrt{\frac{3}{1309}} & 0 & \frac{-3}{z^2}\sqrt{\frac{6}{1309}} & 0 & \frac{-2}{z^3}\sqrt{\frac{3}{1309}} & \frac{-1}{z^3}\sqrt{\frac{6}{715}} \\ \frac{2\sqrt{11}z^4}{25} & \frac{9z^2}{50\sqrt{385}} & \frac{39z}{50\sqrt{385}} & \frac{1497}{250\sqrt{154}} & \frac{57}{250}\sqrt{\frac{21}{374}} & \frac{-36}{125z}\sqrt{\frac{11}{7}} & \frac{-114}{125z}\sqrt{\frac{21}{187}} & \frac{-2\sqrt{11}}{25z^2} & \frac{19}{25z^2}\sqrt{\frac{14}{561}} & \frac{1102}{125\sqrt{2145}z^2} \\ 0 & 12\sqrt{\frac{2}{2387}z^2} & 19\sqrt{\frac{2}{2387}z} & -10\sqrt{\frac{5}{20289}} & 48\sqrt{\frac{3}{20289}} & 0 & \frac{-3}{z}\sqrt{\frac{186}{6545}} & 0 & \frac{-2}{z^2}\sqrt{\frac{93}{6545}} & \frac{14}{5z^2}\sqrt{\frac{6}{4433}} \\ 0 & \frac{z^2}{2}\sqrt{\frac{39}{217}} & \frac{-z}{2}\sqrt{\frac{39}{217}} & \frac{1}{2}\sqrt{\frac{39}{2170}} & \frac{-19}{2}\sqrt{\frac{13}{36890}} & 0 & 0 & 0 & 0 & \frac{-6}{5\sqrt{31}z^2} \\ \frac{2\sqrt{231}z^5}{125} & \frac{3z^3}{1750}\sqrt{\frac{33}{5}} & \frac{54z^2}{875}\sqrt{\frac{3}{55}} & \frac{1149z}{8750}\sqrt{\frac{3}{22}} & \frac{57z}{1250}\sqrt{\frac{2}{374}} & \frac{8541}{8750}\sqrt{\frac{3}{11}} & \frac{741}{625}\sqrt{\frac{187}{187}} & \frac{-18}{125z}\sqrt{\frac{33}{7}} & \frac{-399}{125z}\sqrt{\frac{2}{187}} & \frac{-3306}{625\sqrt{5005}z} \\ 0 & 2\sqrt{\frac{11}{1085}z^3} & \frac{39z^2}{\sqrt{11935}} & 15\sqrt{\frac{2}{2387}z} & \frac{4z}{5}\sqrt{\frac{6}{40579}} & \frac{-25}{\sqrt{2387}} & \frac{142}{5}\sqrt{\frac{3}{40579}} & 0 & \frac{-6}{5z}\sqrt{\frac{186}{1309}} & \frac{-21}{5z}\sqrt{\frac{3}{22165}} \end{array} \right)$$

5.3 Tricritical Ising model: $\mathcal{V}_0 \oplus \mathcal{V}_{3/5}$

Lastly, we present the mode $\phi\left(\frac{3}{5}, 0\right)_{-3/5}$ of the field $\phi\left(\frac{3}{5}, 0\right)(z)$

$$\begin{aligned} \phi\left(\frac{3}{5}, 0\right)_{-3/5} &= (1) \oplus \left(\frac{-4}{5}\sqrt{\frac{3}{385}} \right) \oplus \left(\frac{-3}{25}\sqrt{\frac{7}{11}} \right) \oplus \left(\frac{-263}{875}\sqrt{\frac{3}{22}} \quad \frac{-12}{125}\sqrt{\frac{2}{187}} \right) \oplus \left(\frac{-5177}{4375}\sqrt{\frac{3}{286}} \quad \frac{-372}{625}\sqrt{\frac{2}{2431}} \right) \\ &\oplus \left(25\sqrt{\frac{2}{2387}} \quad \frac{-4}{5}\sqrt{\frac{6}{40579}} \right) \oplus \left(\frac{9}{7}\sqrt{\frac{183}{1426}} \quad -296\sqrt{\frac{2}{739381}} \right) \\ &\oplus \left(\frac{3}{7}\sqrt{\frac{2}{310}} \quad -148\sqrt{\frac{2}{160735}} \right) \oplus \left(0 \quad 5\sqrt{\frac{299}{7259}} \right) \\ &\oplus \left(\frac{-903}{125}\sqrt{\frac{7}{22165}} \quad \frac{-736}{125}\sqrt{\frac{6}{376805}} \quad \frac{-1786}{40625}\sqrt{\frac{3}{2387}} \quad \frac{-1064}{40625}\sqrt{\frac{6}{5797}} \right) \\ &\oplus \left(\frac{125}{\sqrt{41943}} \quad \frac{-63}{25}\sqrt{\frac{14}{237677}} \quad \frac{-42}{25}\sqrt{\frac{14}{908765}} \quad \frac{48}{25}\sqrt{\frac{14}{15449005}} \right) \\ &\oplus \left(9\sqrt{\frac{305}{84847}} \quad \frac{-296}{17}\sqrt{\frac{30}{43493}} \quad \frac{-6}{25}\sqrt{\frac{183}{1103011}} \quad \frac{3304}{1275}\sqrt{\frac{2}{1696227}} \right) \\ &\oplus \left(0 \quad \frac{50}{17}\sqrt{\frac{39}{427}} \quad 0 \quad \frac{-52}{51}\sqrt{\frac{6405}{6405}} \right) \\ &\oplus \left(-\sqrt{\frac{10}{897}} \quad \frac{3}{2}\sqrt{\frac{85}{2093}} \quad \frac{-6}{65}\sqrt{\frac{2}{23}} \quad \frac{36}{65}\sqrt{\frac{7}{391}} \right) \\ &\oplus \left(\frac{1}{7}\sqrt{\frac{34698}{108035}} \quad \frac{33001}{34\sqrt{257256755}} \quad \frac{-592642}{35}\sqrt{\frac{2}{1624392653}} \quad \frac{-80468}{85}\sqrt{\frac{7}{95552509}} \right) \\ &\oplus \left(0 \quad \frac{8}{17}\sqrt{\frac{1495}{121443}} \quad 4\sqrt{\frac{782}{17349}} \quad \frac{407}{17}\sqrt{\frac{23}{121443}} \right) \\ &\oplus \left(\frac{-61541}{4375}\sqrt{\frac{2}{22165}} \quad \frac{-29684}{625}\sqrt{\frac{6}{2637635}} \quad \frac{-36613}{284375}\sqrt{\frac{6}{1705}} \quad \frac{-3116}{40625}\sqrt{\frac{14}{28985}} \right) \\ &\oplus \left(250\sqrt{\frac{5}{964689}} \quad \frac{-16109}{125}\sqrt{\frac{6}{191329985}} \quad \frac{-784}{125}\sqrt{\frac{6}{20901595}} \quad \frac{1792}{125}\sqrt{\frac{21}{355327115}} \right) \\ &\oplus \left(\frac{15}{7}\sqrt{\frac{915}{12121}} \quad \frac{-740}{17}\sqrt{\frac{30}{304451}} \quad \frac{-81}{175}\sqrt{\frac{183}{1575730}} \quad \frac{708}{425}\sqrt{\frac{21}{2827045}} \right) \\ &\oplus \left(0 \quad \frac{125}{119}\sqrt{\frac{39}{61}} \quad 0 \quad \frac{-39}{595}\sqrt{\frac{6}{61}} \right) \\ &\oplus \left(-10\sqrt{\frac{10}{54717}} \quad \frac{15}{2}\sqrt{\frac{255}{127673}} \quad \frac{-116}{65}\sqrt{\frac{3}{7015}} \quad \frac{348}{65}\sqrt{\frac{42}{119255}} \right) \oplus \dots \quad (5.67) \\ &\oplus \left(\frac{2}{7}\sqrt{\frac{57830}{237677}} \quad \frac{33001}{34}\sqrt{\frac{5}{565964861}} \quad \frac{-1185284}{7\sqrt{89341595915}} \quad \frac{-80468}{17}\sqrt{\frac{14}{5255387995}} \right) \\ &\oplus \left(0 \quad \frac{40}{17}\sqrt{\frac{1495}{1335873}} \quad 8\sqrt{\frac{1955}{190839}} \quad \frac{37}{17}\sqrt{\frac{2530}{121443}} \right) \\ &\oplus \left(\frac{-9}{7}\sqrt{\frac{5}{15433}} \quad \frac{64}{9}\sqrt{\frac{34}{540155}} \quad \frac{821}{63}\sqrt{\frac{2}{1003145}} \quad \frac{-8}{3}\sqrt{\frac{7}{17053465}} \right) \\ &\oplus \left(0 \quad \frac{1}{63}\sqrt{\frac{170}{11}} \quad \frac{-4}{9}\sqrt{\frac{70}{143}} \quad \frac{-1}{21}\sqrt{\frac{55}{221}} \right) \end{aligned}$$

$$T(z) = \begin{pmatrix} \frac{1}{10z^2} & \frac{1}{\sqrt{5}z^3} & \frac{1}{2z^4}\sqrt{3} & \frac{2}{z^5}\sqrt{\frac{5}{11}} & -\frac{1}{z^5}\sqrt{\frac{2}{11}} & \frac{25}{4\sqrt{11}z^6} & \frac{-5}{z^6}\sqrt{\frac{10}{341}} & \frac{1}{4z^6}\sqrt{\frac{39}{155}} & \frac{25}{2z^7}\sqrt{\frac{3}{77}} & \frac{-25}{\sqrt{341}z^7} & \frac{3}{2z^7}\sqrt{\frac{39}{1271}} & -\frac{1}{z^7}\sqrt{\frac{3}{287}} \\ \frac{1}{\sqrt{5}z} & \frac{11}{10z^2} & \frac{2}{z^3}\sqrt{\frac{3}{5}} & \frac{13}{2\sqrt{11}z^4} & \frac{1}{z^4}\sqrt{\frac{10}{11}} & \frac{7}{2z^5}\sqrt{\frac{5}{11}} & \frac{19}{z^5}\sqrt{\frac{2}{341}} & \frac{-1}{2z^5}\sqrt{\frac{39}{31}} & \frac{25}{4z^6}\sqrt{\frac{15}{77}} & \frac{15}{z^6}\sqrt{\frac{5}{341}} & \frac{-49}{4z^6}\sqrt{\frac{39}{6355}} & \frac{3}{z^6}\sqrt{\frac{15}{287}} \\ \frac{\sqrt{3}}{2} & \frac{2}{z}\sqrt{\frac{3}{5}} & \frac{21}{10z^2} & \frac{1}{z^3}\sqrt{\frac{33}{5}} & 0 & \frac{23}{4z^4}\sqrt{\frac{3}{11}} & \frac{2}{z^4}\sqrt{\frac{30}{341}} & \frac{1}{4z^4}\sqrt{\frac{65}{31}} & \frac{30}{\sqrt{77}z^5} & \frac{13}{z^5}\sqrt{\frac{3}{341}} & \frac{2}{z^5}\sqrt{\frac{13}{1271}} & \frac{-15}{\sqrt{287}z^5} \\ 2\sqrt{\frac{5}{11}}z & \frac{13}{2\sqrt{11}} & \frac{1}{z}\sqrt{\frac{33}{5}} & \frac{31}{10z^2} & 0 & \frac{8}{\sqrt{5}z^3} & 0 & 0 & \frac{11}{4z^4}\sqrt{\frac{15}{7}} & \frac{1}{z^4}\sqrt{\frac{5}{31}} & \frac{1}{4z^4}\sqrt{\frac{2145}{1271}} & \frac{5}{z^4}\sqrt{\frac{15}{3157}} \\ -\sqrt{\frac{2}{11}}z & \sqrt{\frac{10}{11}} & 0 & 0 & \frac{31}{10z^2} & 0 & \frac{1}{z^3}\sqrt{\frac{31}{5}} & 0 & 0 & \frac{1}{2z^4}\sqrt{\frac{31}{2}} & 0 & \frac{1}{2z^4}\sqrt{\frac{861}{22}} \\ \frac{25z^2}{4\sqrt{11}} & \frac{7z}{2}\sqrt{\frac{5}{11}} & \frac{23}{4}\sqrt{\frac{3}{11}} & \frac{8}{\sqrt{5}z} & 0 & \frac{41}{10z^2} & 0 & 0 & \frac{\sqrt{21}}{z^3} & 0 & 0 & 0 \\ -5\sqrt{\frac{10}{341}}z^2 & 19\sqrt{\frac{2}{341}}z & 2\sqrt{\frac{30}{341}} & 0 & \frac{1}{z}\sqrt{\frac{31}{5}} & 0 & \frac{41}{10z^2} & 0 & 0 & \frac{6}{z^3}\sqrt{\frac{2}{5}} & 0 & 0 \\ \frac{1z^2}{4}\sqrt{\frac{39}{155}} & \frac{1z}{2}\sqrt{\frac{39}{31}} & \frac{1}{4}\sqrt{\frac{65}{31}} & 0 & 0 & 0 & 0 & \frac{41}{10z^2} & 0 & 0 & \frac{1}{z^3}\sqrt{\frac{41}{5}} & 0 \\ \frac{25z^3}{2}\sqrt{\frac{3}{77}} & \frac{25z^2}{4}\sqrt{\frac{15}{77}} & \frac{30z}{\sqrt{77}} & \frac{11}{4}\sqrt{\frac{15}{7}} & 0 & \frac{\sqrt{21}}{z} & 0 & 0 & \frac{51}{10z^2} & 0 & 0 & 0 \\ \frac{-25z^3}{\sqrt{341}} & 15\sqrt{\frac{5}{341}}z^2 & 13\sqrt{\frac{3}{341}}z & \sqrt{\frac{5}{31}} & \frac{1}{2}\sqrt{\frac{31}{2}} & 0 & \frac{6}{z}\sqrt{\frac{2}{5}} & 0 & 0 & \frac{51}{10z^2} & 0 & 0 \\ \frac{3z^3}{2}\sqrt{\frac{39}{1271}} & \frac{-49z^2}{4}\sqrt{\frac{39}{6355}} & 2\sqrt{\frac{13}{1271}}z & \frac{1}{4}\sqrt{\frac{2145}{1271}} & 0 & 0 & 0 & \frac{1}{z}\sqrt{\frac{41}{5}} & 0 & 0 & \frac{51}{10z^2} & 0 \\ -\sqrt{\frac{3}{287}}z^3 & 3\sqrt{\frac{15}{287}}z^2 & \frac{-15z}{\sqrt{287}} & 5\sqrt{\frac{15}{3157}} & \frac{1}{2}\sqrt{\frac{861}{22}} & 0 & 0 & 0 & 0 & 0 & 0 & \frac{51}{10z^2} \end{pmatrix}$$

Figure 9: Energy-momentum tensor (in the L_1 basis) of the tricritical Ising model in the $h = 1/10$ sector.

$$\phi\left(\begin{matrix} 3/5 \\ 3/5, 0 \end{matrix}\right)(z) =$$

1	$\frac{6}{\sqrt{35}z^2}$	$\frac{6}{\sqrt{35}z^3}$	$\frac{9}{5z^4}\sqrt{\frac{2}{7}}$	$\frac{8}{5z^4}\sqrt{\frac{6}{119}}$	$\frac{12}{5\sqrt{7}z^5}$	$\frac{16}{5z^5}\sqrt{\frac{3}{119}}$	$\frac{6}{7z^6}$	$\frac{8}{z^6}\sqrt{\frac{2}{357}}$	$\frac{47}{35\sqrt{195}z^6}$	$\frac{4}{5z^6}\sqrt{\frac{14}{3315}}$
$\sqrt{\frac{6}{5}z}$	$\frac{\sqrt{6}}{5z^2}$	$\frac{\sqrt{6}}{5z^2}$	$\frac{8\sqrt{\frac{3}{35}}}{5z^2}$	$\frac{-16\sqrt{\frac{7}{85}}}{5z^2}$	$\frac{\sqrt{42}}{5z^2}$	$\frac{-72\sqrt{\frac{595}{5z^2}}}{5z^2}$	$\frac{4\sqrt{\frac{6}{7z^2}}}{7z^2}$	$\frac{-16}{\sqrt{595}z^5}$	$\frac{-188\sqrt{\frac{13}{221}}}{25z^5}$	$\frac{-32\sqrt{\frac{7}{221}}}{25z^5}$
$\frac{\sqrt{33}z^2}{5}$	$\frac{-4\sqrt{\frac{3}{385}}}{5}$	$\frac{-4\sqrt{\frac{3}{385}}}{5}$	$\frac{-\sqrt{\frac{66}{25z^2}}}{25z^2}$	$\frac{72\sqrt{\frac{14}{187}}}{25z^2}$	$\frac{2\sqrt{\frac{39}{25z^3}}}{25z^3}$	$\frac{-16\sqrt{\frac{7}{187}}}{25z^3}$	$\frac{2\sqrt{\frac{33}{35z^4}}}{35z^4}$	$\frac{-64\sqrt{\frac{9}{1309}}}{5z^4}$	$\frac{3572}{175\sqrt{715}z^4}$	$\frac{304\sqrt{\frac{14}{12155}}}{25z^4}$
0	$\frac{\sqrt{\frac{65}{77}}}{\sqrt{77}}$	0	0	$\frac{2\sqrt{\frac{78}{1309}}}{z^2}$	0	$\frac{2\sqrt{\frac{39}{1309}}}{z^3}$	0	$\frac{\sqrt{\frac{78}{1309}}}{z^4}$	$\frac{2\sqrt{\frac{3}{55}}}{z^4}$	$\frac{-16\sqrt{\frac{6}{6545}}}{z^4}$
$4\sqrt{\frac{11}{5}z^3}$	$\frac{-6z}{25\sqrt{77}}$	$\frac{-3\sqrt{\frac{7}{11}}}{25}$	$\frac{-24\sqrt{\frac{22}{35}}}{25z}$	$\frac{-24\sqrt{\frac{42}{935}}}{25z}$	$\frac{-12\sqrt{\frac{11}{35}}}{25z^2}$	$\frac{12\sqrt{\frac{231}{85}}}{25z^2}$	0	$\frac{4\sqrt{\frac{14}{2805}}}{z^3}$	$\frac{-1786}{125\sqrt{429}z^3}$	$\frac{-1064\sqrt{\frac{14}{7293}}}{125z^3}$
0	$3\sqrt{\frac{2}{77}z}$	$5\sqrt{\frac{2}{77}}$	0	$\frac{-8\sqrt{\frac{3}{6545}}}{z}$	0	$\frac{-504}{125\sqrt{187}z}$	0	$\frac{16\sqrt{\frac{3}{6545}}}{z^3}$	$\frac{-14\sqrt{\frac{6}{143}}}{5z^3}$	$\frac{32\sqrt{\frac{21}{5z^3}}}{5z^3}$
$\frac{2\sqrt{231}z^4}{25}$	$\frac{-\sqrt{\frac{35}{175}}z^2}{175}$	$\frac{-31\sqrt{\frac{3}{55}z}}{175}$	$\frac{-263\sqrt{\frac{3}{875}}}{875}$	$\frac{-12\sqrt{\frac{2}{187}}}{125}$	$\frac{-16\sqrt{\frac{33}{125z}}}{125z}$	$\frac{-504}{125\sqrt{187}z}$	$\frac{-4\sqrt{\frac{33}{25z^2}}}{25z^2}$	$\frac{168\sqrt{\frac{2}{187}}}{25z^2}$	$\frac{2679}{125\sqrt{5005}z^2}$	$\frac{1596\sqrt{\frac{2}{12155}}}{125z^2}$
0	$3\sqrt{\frac{183}{31}z^2}$	$\frac{-3\sqrt{\frac{183}{31}z}}{7}$	$25\sqrt{\frac{2}{2387}}$	$\frac{-4\sqrt{\frac{6}{40579}}}{5}$	0	$\frac{-4\sqrt{\frac{93}{1309}}}{5z}$	0	$\frac{3\sqrt{\frac{186}{1309}}}{5z^2}$	$\frac{126\sqrt{\frac{3}{22165}}}{5z^2}$	$\frac{-144\sqrt{\frac{42}{376805}}}{5z^2}$
0	$\frac{24\sqrt{\frac{366}{3565}z^3}}{7}$	$\frac{-3\sqrt{\frac{183}{713}z^2}}{7}$	$\frac{-66\sqrt{\frac{183}{713}}}{35}$	$\frac{-888z}{5\sqrt{739381}}$	0	$\frac{-296\sqrt{\frac{2}{739381}}}{z}$	0	0	$\frac{9\sqrt{\frac{61}{2821}}}{5z^2}$	$\frac{-1652\sqrt{\frac{2}{417911}}}{15z^2}$
0	0	0	0	$5\sqrt{\frac{299}{7259}}$	0	$10\sqrt{\frac{65}{7259}}$	0	0	0	$\frac{26\sqrt{\frac{115}{7239}}}{3z^2}$
$\frac{2\sqrt{6006}z^5}{125}$	$\frac{-8\sqrt{\frac{66}{875}z^3}}{875}$	$\frac{-41\sqrt{\frac{33}{130}z^2}}{875}$	$\frac{-1564\sqrt{\frac{3}{4375}}}{4375}$	$\frac{-72z}{625\sqrt{2431}}$	$\frac{-5177\sqrt{\frac{3}{286}}}{4375}$	$\frac{-372\sqrt{\frac{2}{2431}}}{625}$	$\frac{-8\sqrt{\frac{85}{125z}}}{125z}$	$\frac{-108\sqrt{\frac{13}{187}}}{125z}$	$\frac{-5358\sqrt{\frac{2}{8125z}}}{8125z}$	$\frac{-6384}{8125\sqrt{935}z}$
0	$\frac{8\sqrt{\frac{11}{651}z^3}}{5}$	$\frac{\sqrt{\frac{33}{217}z^2}}{217}$	$5\sqrt{\frac{30}{2387}z}$	$\frac{-4\sqrt{\frac{2}{202805}z}}{5}$	$25\sqrt{\frac{5}{7161}}$	$\frac{-2\sqrt{\frac{17}{11935}}}{5}$	0	$\frac{-12\sqrt{\frac{6545}{5z}}}{5z}$	$\frac{-84}{25\sqrt{4433z}}$	$\frac{96\sqrt{\frac{14}{75361}}}{25z}$
0	$\frac{24\sqrt{\frac{366}{3565}z^3}}{7}$	$\frac{-3\sqrt{\frac{183}{713}z^2}}{7}$	$\frac{-66\sqrt{\frac{183}{713}}}{35}$	$\frac{-888z}{5\sqrt{739381}}$	$\frac{9\sqrt{\frac{183}{1426}}}{7}$	$\frac{-296\sqrt{\frac{2}{739381}}}{z}$	0	0	$\frac{-18\sqrt{\frac{192}{324415}}}{5z}$	$\frac{6608}{15\sqrt{48059765}z}$
0	0	0	0	$3\sqrt{\frac{130}{7259}z}$	0	$10\sqrt{\frac{65}{7259}}$	0	0	0	$\frac{-52\sqrt{\frac{2}{7259}}}{3z}$
0	$4\sqrt{\frac{5}{897}z^3}$	$\frac{-2\sqrt{\frac{15}{299}z^2}}{299}$	$2\sqrt{\frac{6}{299}z}$	$\frac{-2\sqrt{\frac{34}{299}z}}{299}$	$\frac{-2}{\sqrt{897}}$	$\sqrt{\frac{17}{299}}$	0	0	$\frac{-24\sqrt{\frac{7}{115}}}{13z}$	$\frac{504\sqrt{\frac{2}{1955}}}{13z}$
$\frac{2\sqrt{31031}z^6}{125}$	$\frac{-24\sqrt{\frac{11}{403}z^4}}{625}$	$\frac{-408\sqrt{\frac{11}{403}z^3}}{4375}$	$\frac{-2589\sqrt{\frac{1030}{4375}z^2}}{4375}$	$\frac{-12\sqrt{\frac{66}{34255}z^2}}{625}$	$\frac{-38991z}{4375\sqrt{22165}}$	$\frac{-984\sqrt{\frac{3}{376805}z}}{625}$	$\frac{-903\sqrt{\frac{7}{22165}}}{125}$	$\frac{-736\sqrt{\frac{6}{376805}}}{125}$	$\frac{-1786\sqrt{\frac{3}{2387}}}{40625}$	$\frac{-1064\sqrt{\frac{6}{5797}}}{40625}$
0	$\frac{4\sqrt{\frac{231}{6355}z^4}}{5}$	$\frac{32\sqrt{\frac{11}{133455}z^3}}{35}$	$5\sqrt{\frac{66}{8897}z^2}$	$\frac{-2\sqrt{\frac{22}{151249}z^2}}{25}$	$100\sqrt{\frac{3}{97867}z}$	$\frac{-142z}{25\sqrt{1663739}}$	$\frac{125}{\sqrt{41943}}$	$\frac{-63\sqrt{\frac{14}{237677}}}{25}$	$\frac{-42}{25\sqrt{908765}}$	$\frac{48\sqrt{\frac{14}{15449005}}}{25}$
0	$\frac{72\sqrt{\frac{61}{12121}z^4}}{5}$	$\frac{216\sqrt{\frac{61}{12121}z^3}}{35}$	$\frac{-801\sqrt{\frac{61}{12121}z^2}}{35}$	$\frac{-1628\sqrt{\frac{6}{217465}z^2}}{85}$	$\frac{-171\sqrt{\frac{61}{60605}z}}{7}$	$\frac{-1184\sqrt{\frac{3}{217465}z}}{17}$	$9\sqrt{\frac{305}{84847}}$	$\frac{-296\sqrt{\frac{30}{43493}}}{17}$	$\frac{-6\sqrt{\frac{183}{1103011}}}{25}$	$\frac{3304\sqrt{\frac{2}{1696227}}}{1275}$
0	0	0	0	$\frac{11\sqrt{\frac{39}{427}z^2}}{17}$	0	$20\sqrt{\frac{78}{427}z}$	0	$\frac{50\sqrt{\frac{39}{427}}}{17}$	0	$\frac{-52}{51\sqrt{6405}}$
0	$\sqrt{\frac{42}{299}z^4}$	$\frac{-2\sqrt{\frac{14}{897}z^3}}{897}$	$\frac{-2\sqrt{\frac{21}{1495}z^2}}{1495}$	$\frac{-11\sqrt{\frac{17}{10465}z^2}}{10465}$	$2\sqrt{\frac{42}{1495}z}$	$\frac{-17\sqrt{\frac{17}{20930}z}}{20930}$	$\frac{-\sqrt{\frac{10}{897}}}{897}$	$\frac{3\sqrt{\frac{85}{2093}}}{2}$	$\frac{-6\sqrt{\frac{2}{25}}}{65}$	$\frac{36\sqrt{\frac{7}{391}}}{65}$
0	$\sqrt{\frac{34698}{151249}z^4}$	$\frac{-2\sqrt{\frac{34698}{151249}z^3}}{151249}$	$\frac{-2\sqrt{\frac{17349}{756245}z^2}}{756245}$	$\frac{99003z^2}{17\sqrt{257256755}}$	$\frac{-2\sqrt{\frac{34698}{756245}z}}{756245}$	$\frac{-99003z}{17\sqrt{514513510}}$	$\frac{\sqrt{\frac{34698}{108035}}}{7}$	$\frac{33001}{34\sqrt{257256755}}$	$\frac{-592642\sqrt{\frac{2}{1624392653}}}{35}$	$\frac{-80468\sqrt{\frac{7}{9552509}}}{85}$
0	0	0	0	$\frac{16\sqrt{\frac{4485}{40481}z^2}}{17\sqrt{40481}}$	0	$\frac{-8\sqrt{\frac{8970}{40481}z}}{17\sqrt{40481}}$	0	$\frac{8\sqrt{\frac{1495}{17\sqrt{121443}}}}{17\sqrt{121443}}$	$4\sqrt{\frac{782}{17349}}$	$\frac{407\sqrt{\frac{23}{121443}}}{17\sqrt{121443}}$

Figure 11: Primary field $\phi\left(\begin{matrix} 3/5 \\ 3/5, 0 \end{matrix}\right)(z)$ (in the L_1 basis) of the tricritical Ising model. The squares of the first column entries agree with $(1-u)^{-6/5}$.

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